

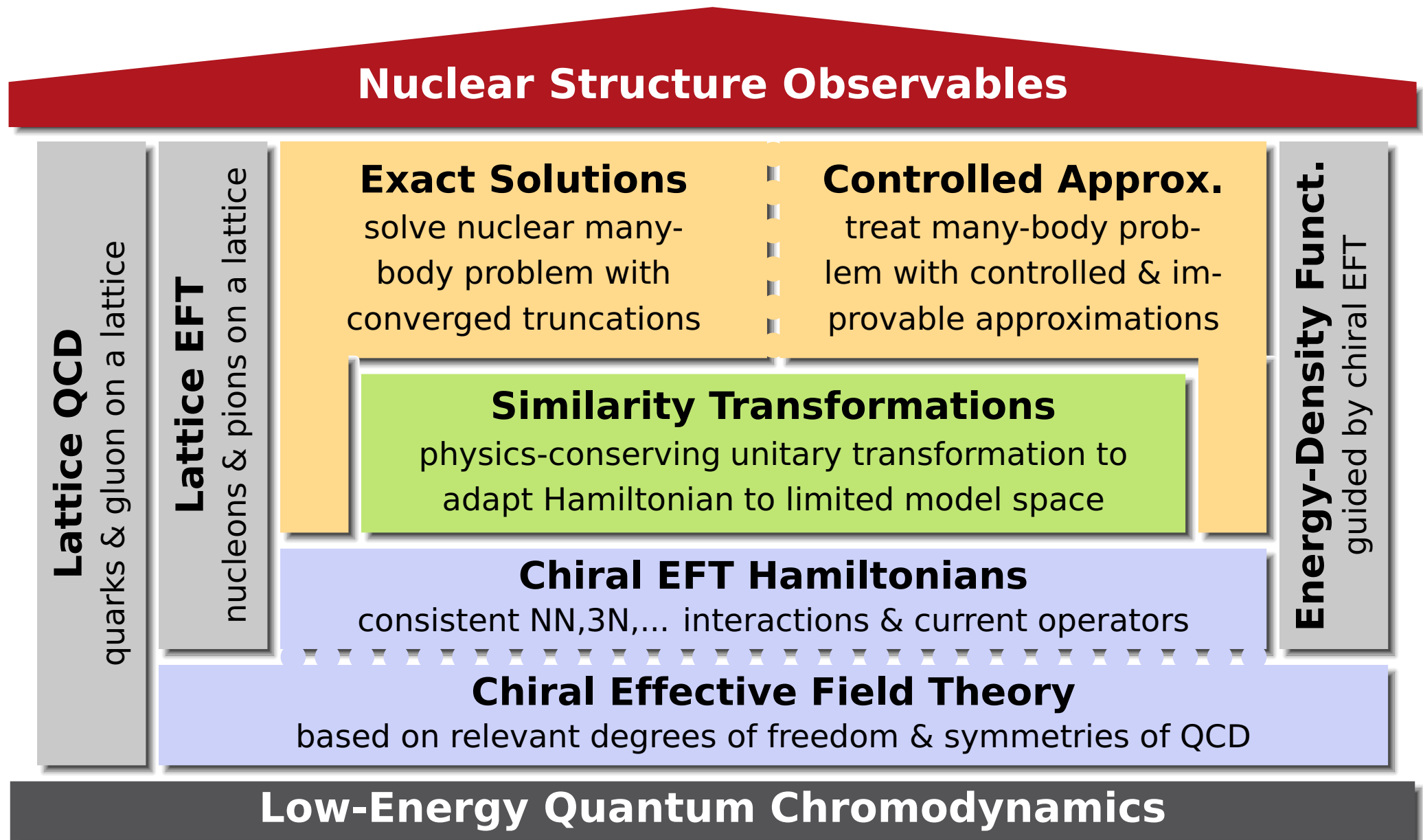
New Horizons in Ab Initio Nuclear Structure Theory

Robert Roth

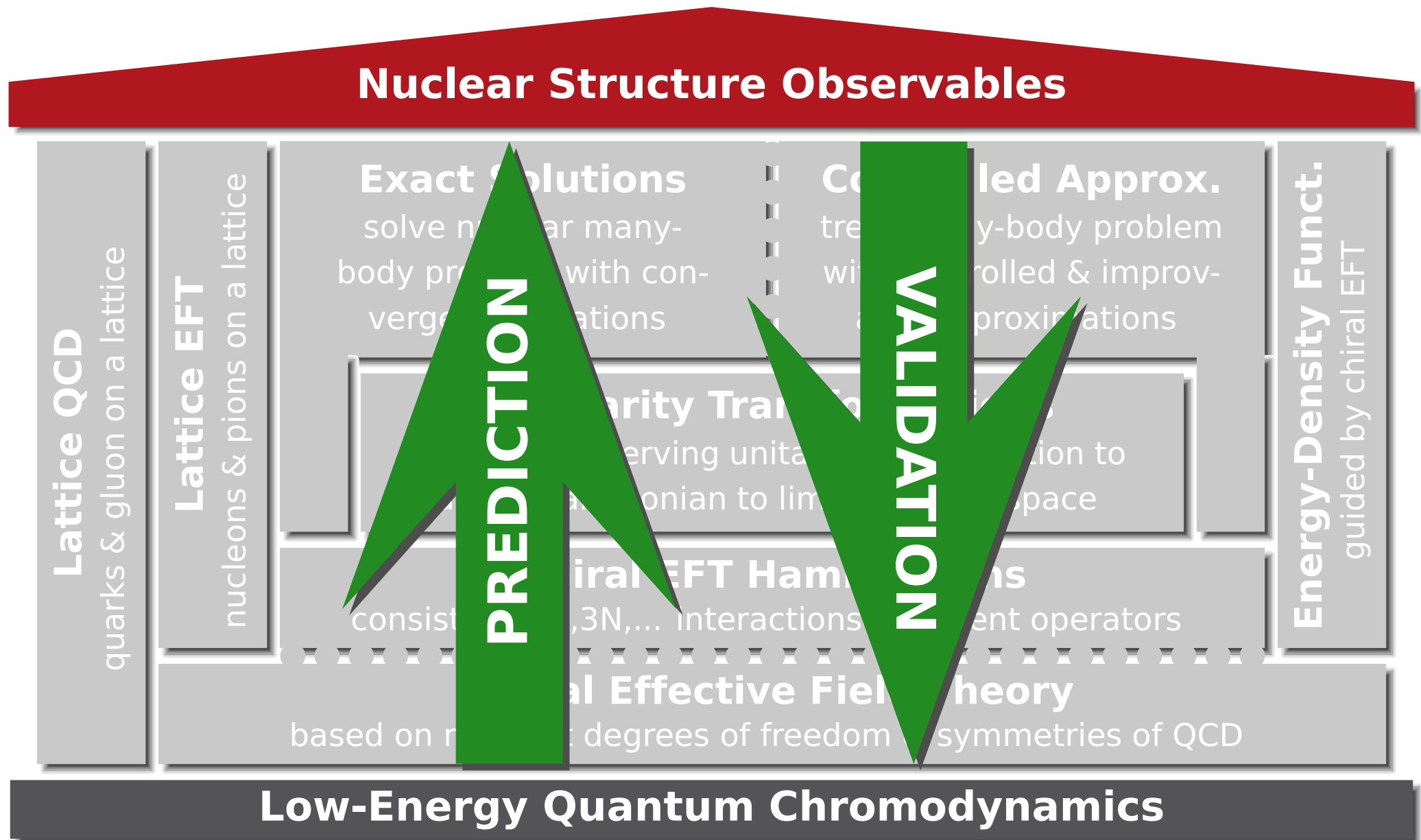


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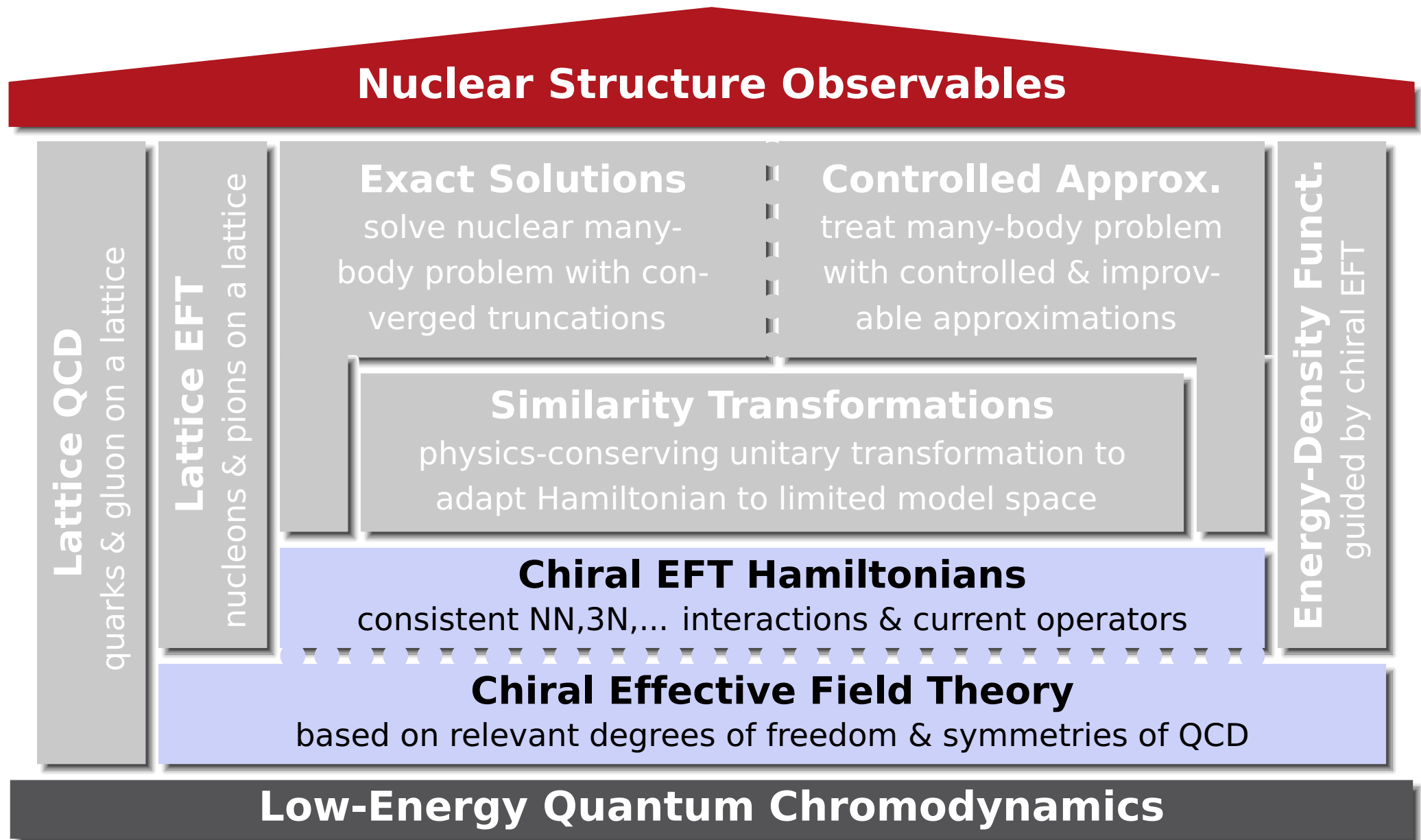
Ab Initio Nuclear Structure



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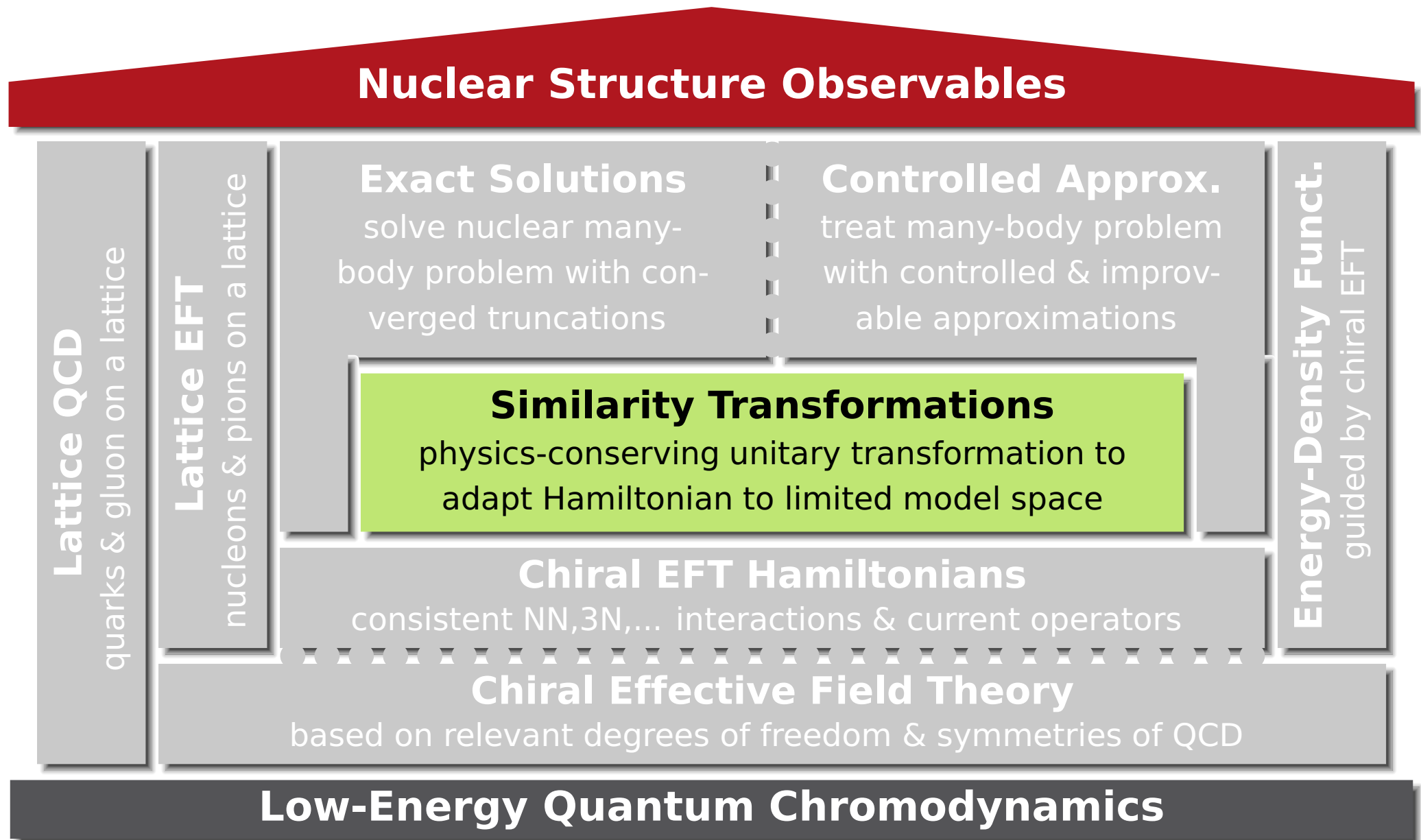
Nuclear Interactions from Chiral EFT

Weinberg, van Kolck, Machleidt, Entem, Meissner, Epelbaum, Krebs, Bernard,...

- **chiral EFT background:** talks by Machleidt & Epelbaum
- **standard Hamiltonian:**
 - NN at N3LO: Entem & Machleidt, 500 MeV cutoff
 - 3N at N2LO: Navrátil, A=3 fit, 500 MeV cutoff
- **alternatives:**
 - modified 3N interaction at N2LO (cutoff, LECs)
 - consistent Hamiltonians at N2LO (NN: Epelbaum, POUNDERS-opt.)
 - consistent Hamiltonians at N3LO
 - Δ -full chiral EFT, YN interaction,...

	NN	3N	4N
LO			
NLO			
N ² LO			
N ³ LO			
	+ ...	+ ...	+ ...

Ab Initio Nuclear Structure



Similarity Renormalization Group

Wegner, Glazek, Wilson, Perry, Bogner, Furnstahl, Hergert, Roth, Jurgenson, Navratil,...

continuous transformation driving
Hamiltonian to band-diagonal form
with respect to a chosen basis

- **unitary transformation** of Hamiltonian

$$H_\alpha = U_\alpha^\dagger H U_\alpha$$

simplicity and flexibility
are great advantages of
the SRG approach

- **evolution equations** for H_α and U_α

$$\frac{d}{d\alpha} H_\alpha = [\eta_\alpha, H_\alpha]$$

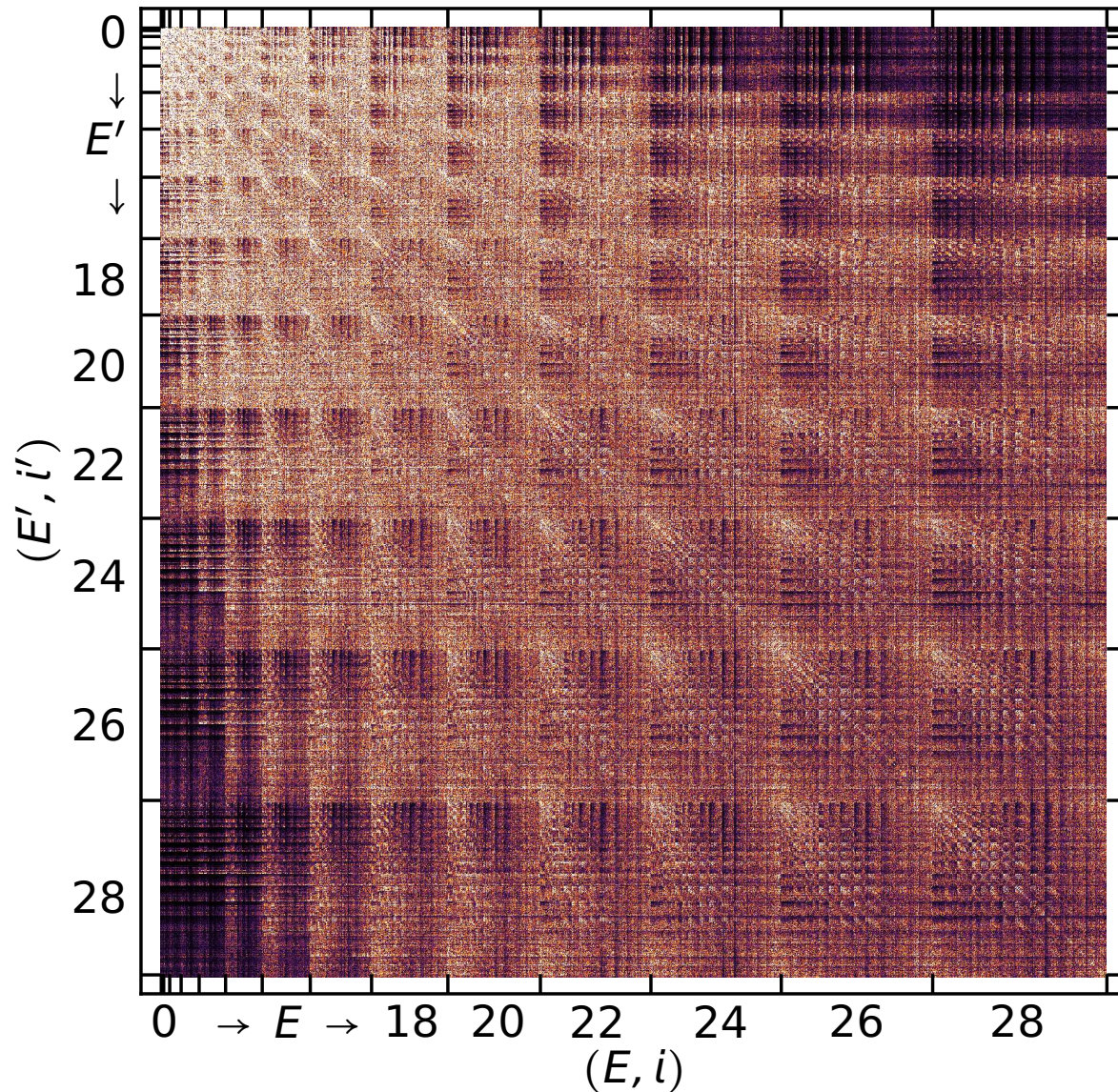
solve SRG evolution
equations using two-,
three- & four-body matrix
representation

- **dynamic generator**: commutator with the operator in whose
eigenbasis H_α shall be diagonalized

$$\eta_\alpha = (2\mu)^2 [T_{\text{int}}, H_\alpha]$$

SRG Evolution in Three-Body Space

3B-Jacobi HO matrix elements

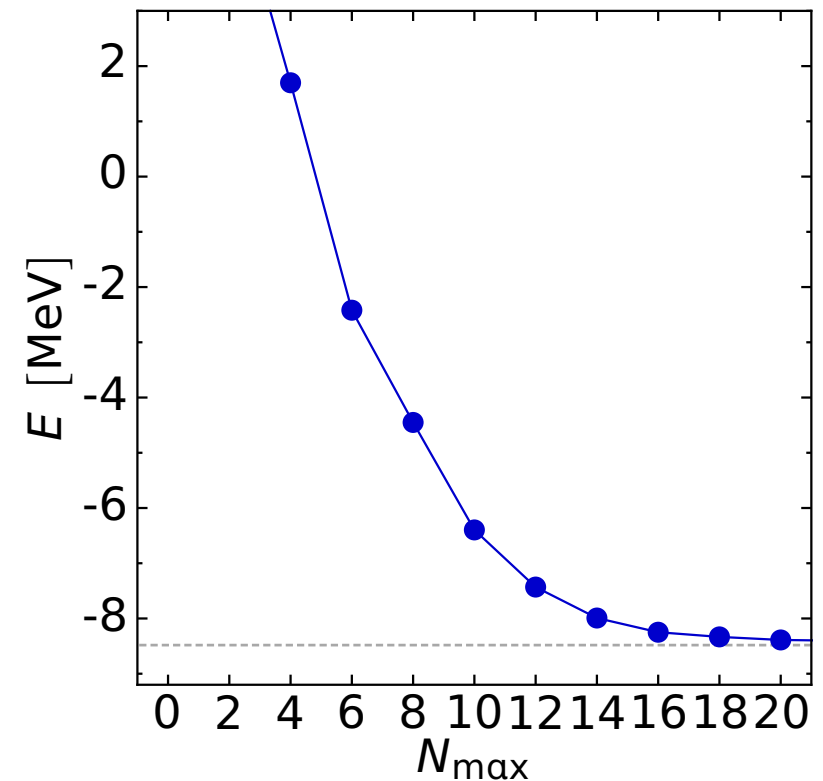


$$\alpha = 0.000 \text{ fm}^4$$

$$\Lambda = \infty \text{ fm}^{-1}$$

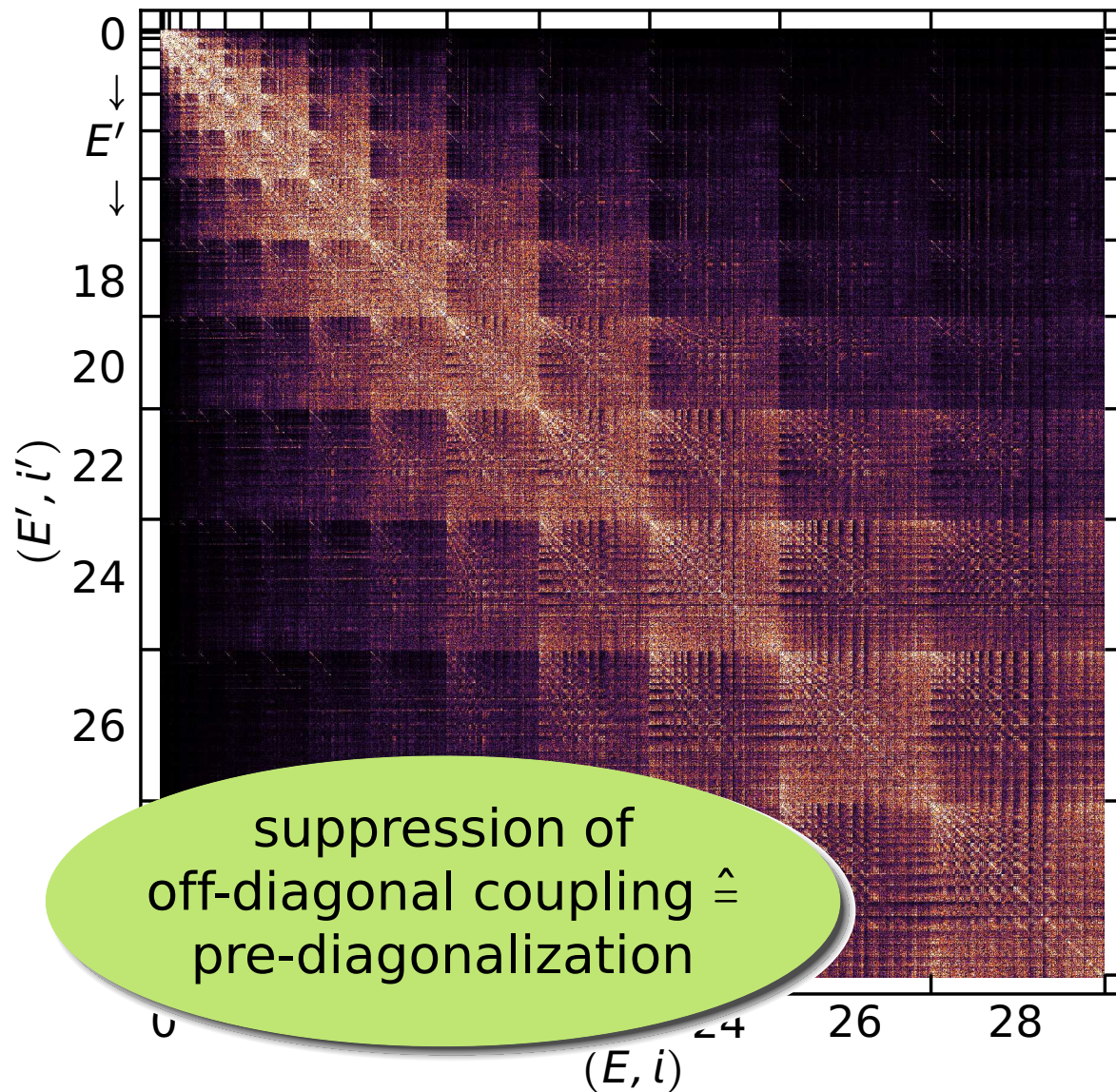
$$J^\pi = \frac{1}{2}^+, T = \frac{1}{2}, \hbar\Omega = 28 \text{ MeV}$$

NCSM ground state ${}^3\text{H}$



SRG Evolution in Three-Body Space

3B-Jacobi HO matrix elements

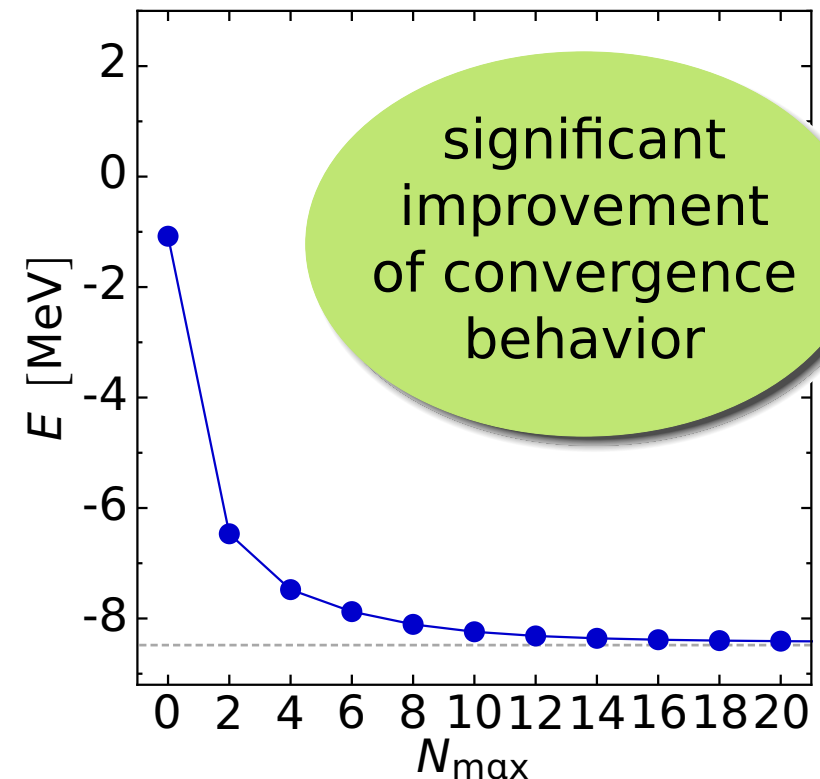


$$\alpha = 0.320 \text{ fm}^4$$

$$\Lambda = 1.33 \text{ fm}^{-1}$$

$$J^\pi = \frac{1}{2}^+, T = \frac{1}{2}, \hbar\Omega = 28 \text{ MeV}$$

NCSM ground state ^3H



Hamiltonian in A-Body Space

- evolution **induces n -body contributions** $H_{\alpha}^{[n]}$ to Hamiltonian

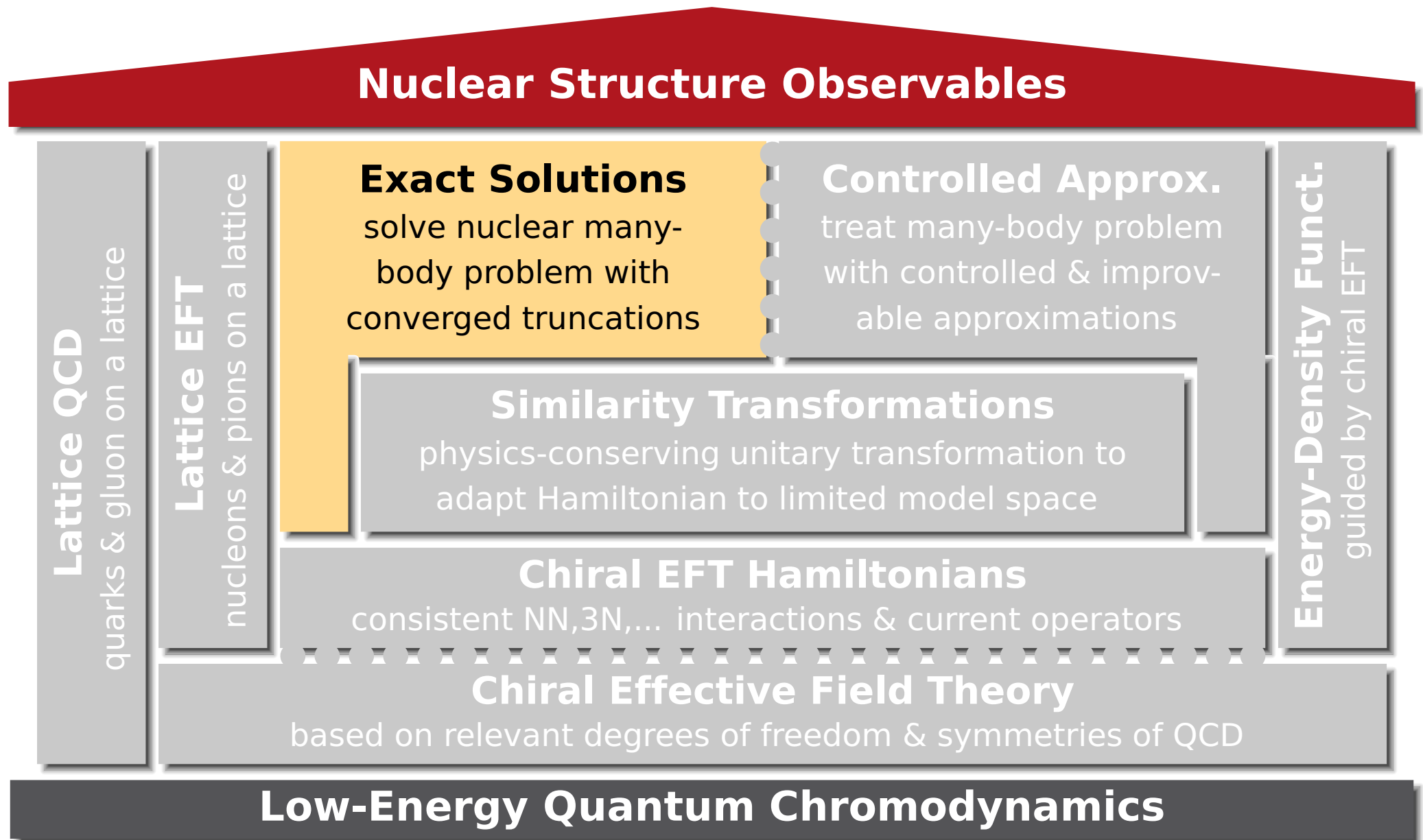
$$H_{\alpha} = H_{\alpha}^{[1]} + H_{\alpha}^{[2]} + H_{\alpha}^{[3]} + H_{\alpha}^{[4]} + \dots$$

- **truncation of cluster series** formally destroys unitarity and invariance of energy eigenvalues (independence of α)
- flow-parameter provides **diagnostic tool** to assess neglected higher-order contributions

SRG-Evolved Hamiltonians

NN_{only}	use initial NN, keep evolved NN
NN + 3N_{ind}	use initial NN, keep evolved NN+3N
NN + 3N_{full}	use initial NN+3N, keep evolved NN+3N
NN + 3N_{full} + 4N_{ind}	use initial NN+3N, keep evolved NN+3N+4N

Ab Initio Nuclear Structure



No-Core Shell Model

Barrett, Vary, Navratil, Maris, Nogga, Roth,...

NCSM is one of the most powerful and universal exact ab-initio methods

- construct matrix representation of Hamiltonian using a **basis of HO Slater determinants** truncated w.r.t. HO excitation energy $N_{\max} \hbar \Omega$
- solve **large-scale eigenvalue problem** for a few extremal eigenvalues
- **all relevant observables** can be computed from the eigenstates
- range of applicability limited by **factorial growth** of basis with $N_{\max}$ & A
- adaptive **importance truncation** extends the range of NCSM by reducing the model space to physically relevant states

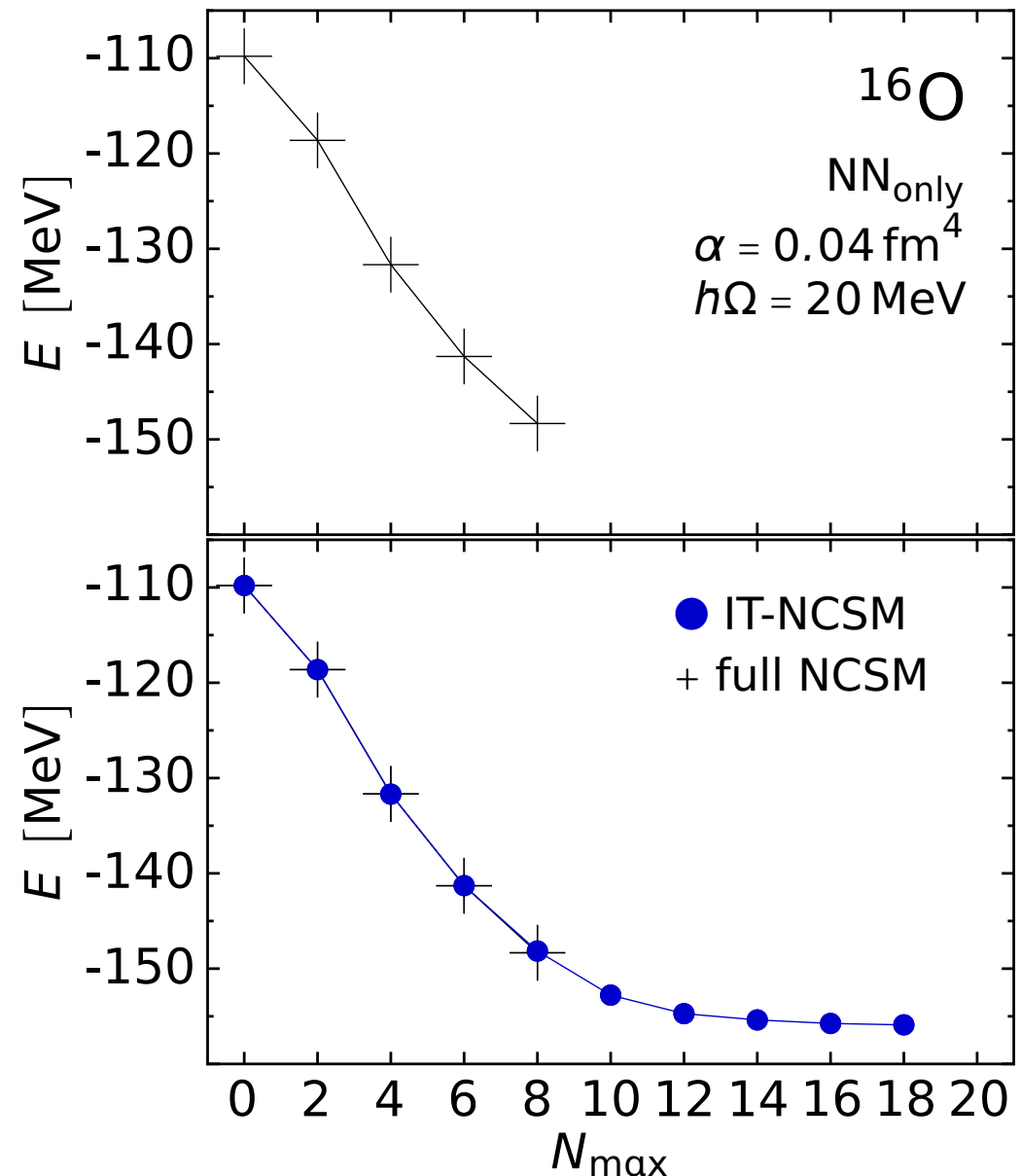
Importance Truncated NCSM

Roth, PRC 79, 064324 (2009); PRL 99, 092501 (2007)

- converged NCSM calculations essentially restricted to lower/mid p-shell
- full $N_{\max} = 10$ calculation for ^{16}O very difficult (basis dimension $> 10^{10}$)

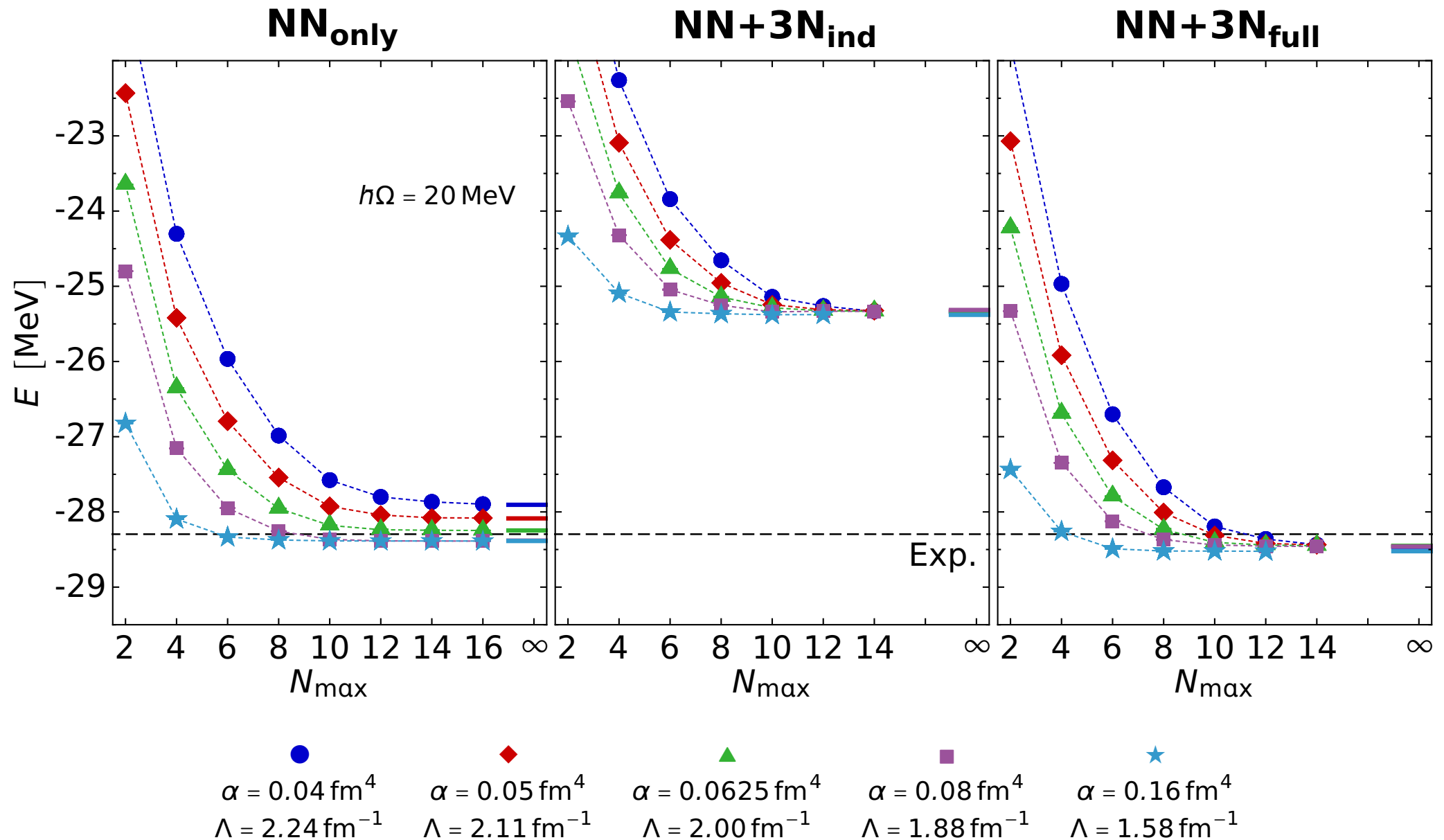
Importance Truncation

reduce model space to the relevant basis states using an **a priori importance measure** derived from MBPT



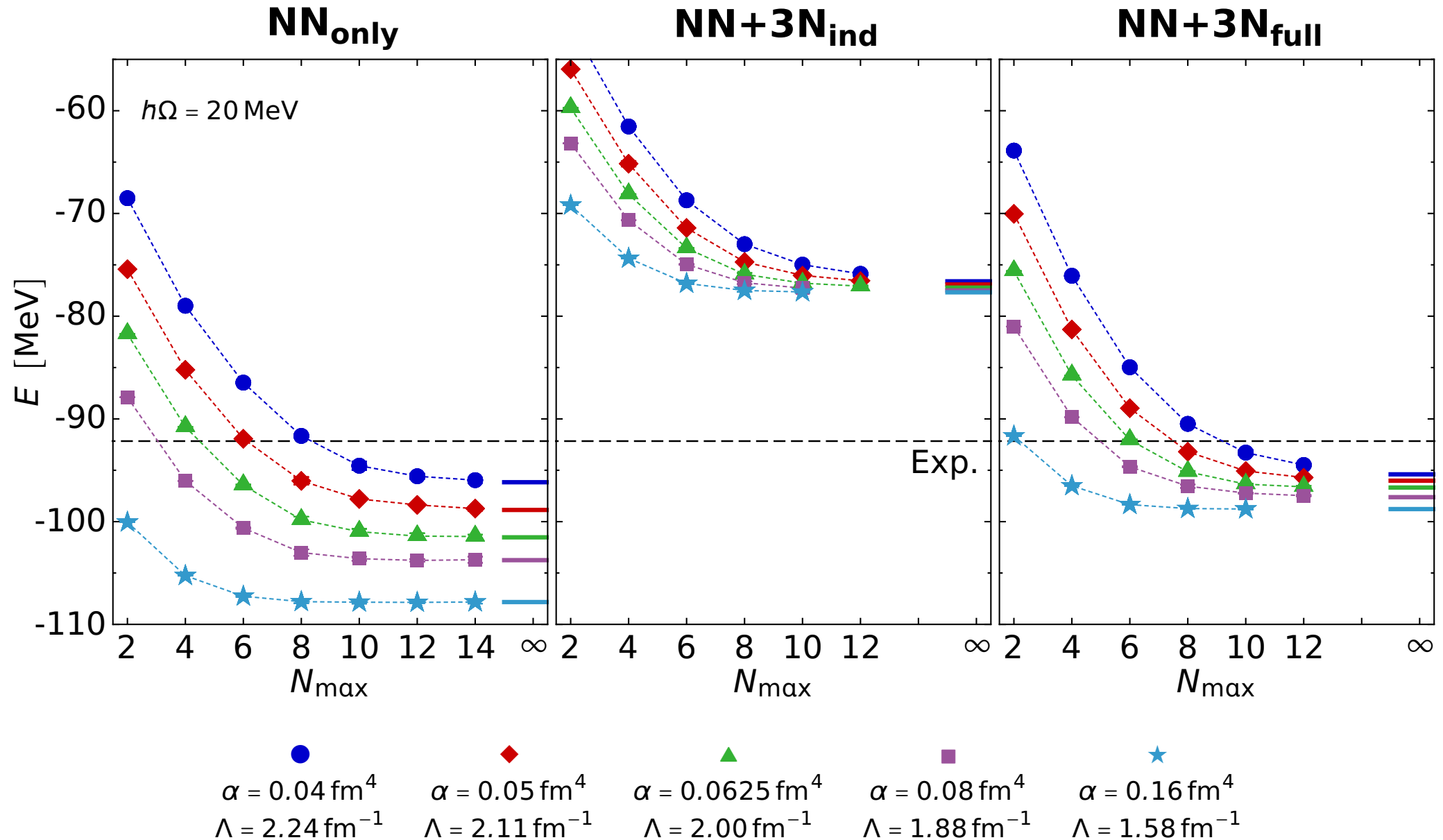
^4He : Ground-State Energies

Roth, et al; PRL 107, 072501 (2011)



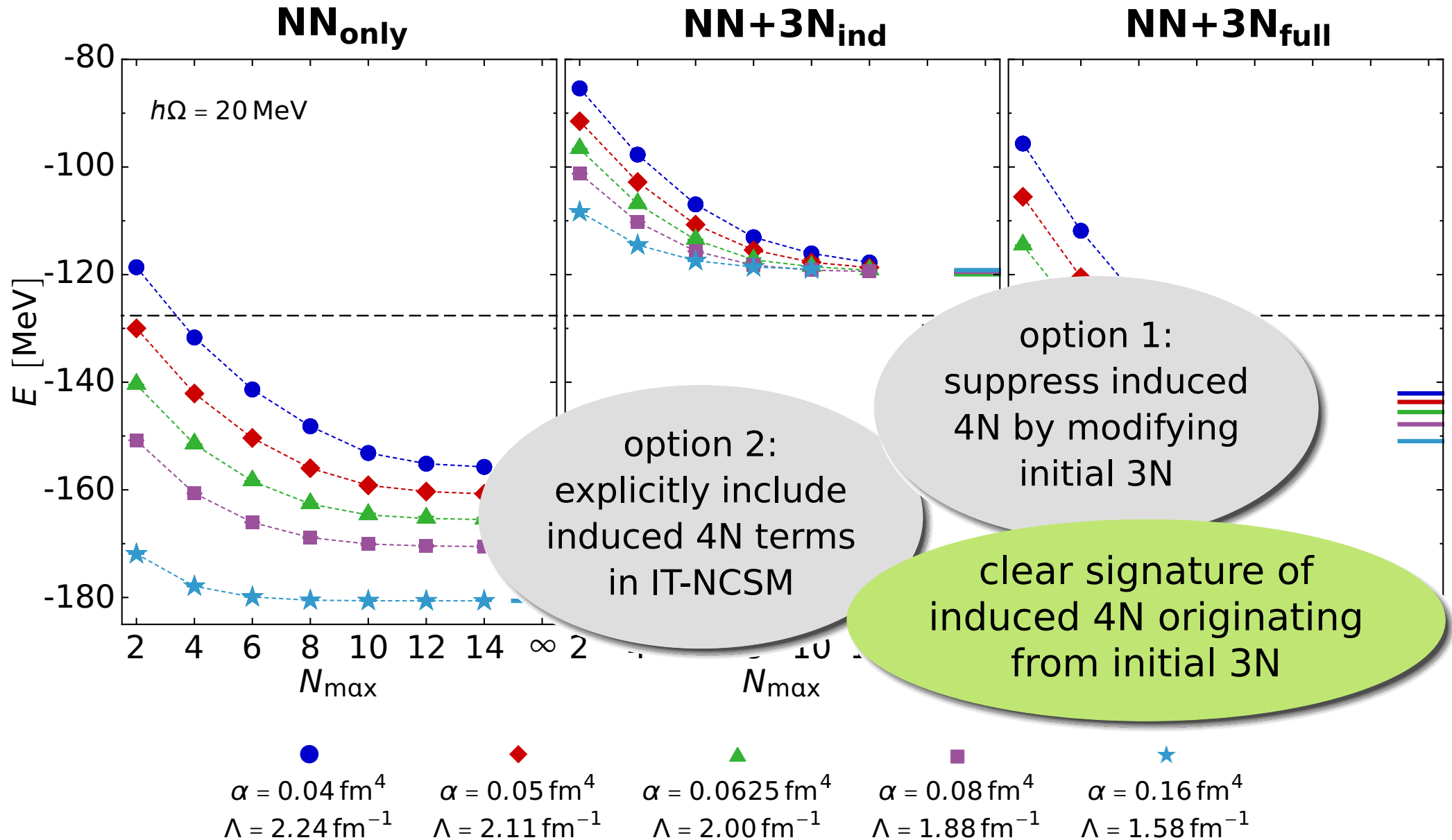
^{12}C : Ground-State Energies

Roth, et al; PRL 107, 072501 (2011)

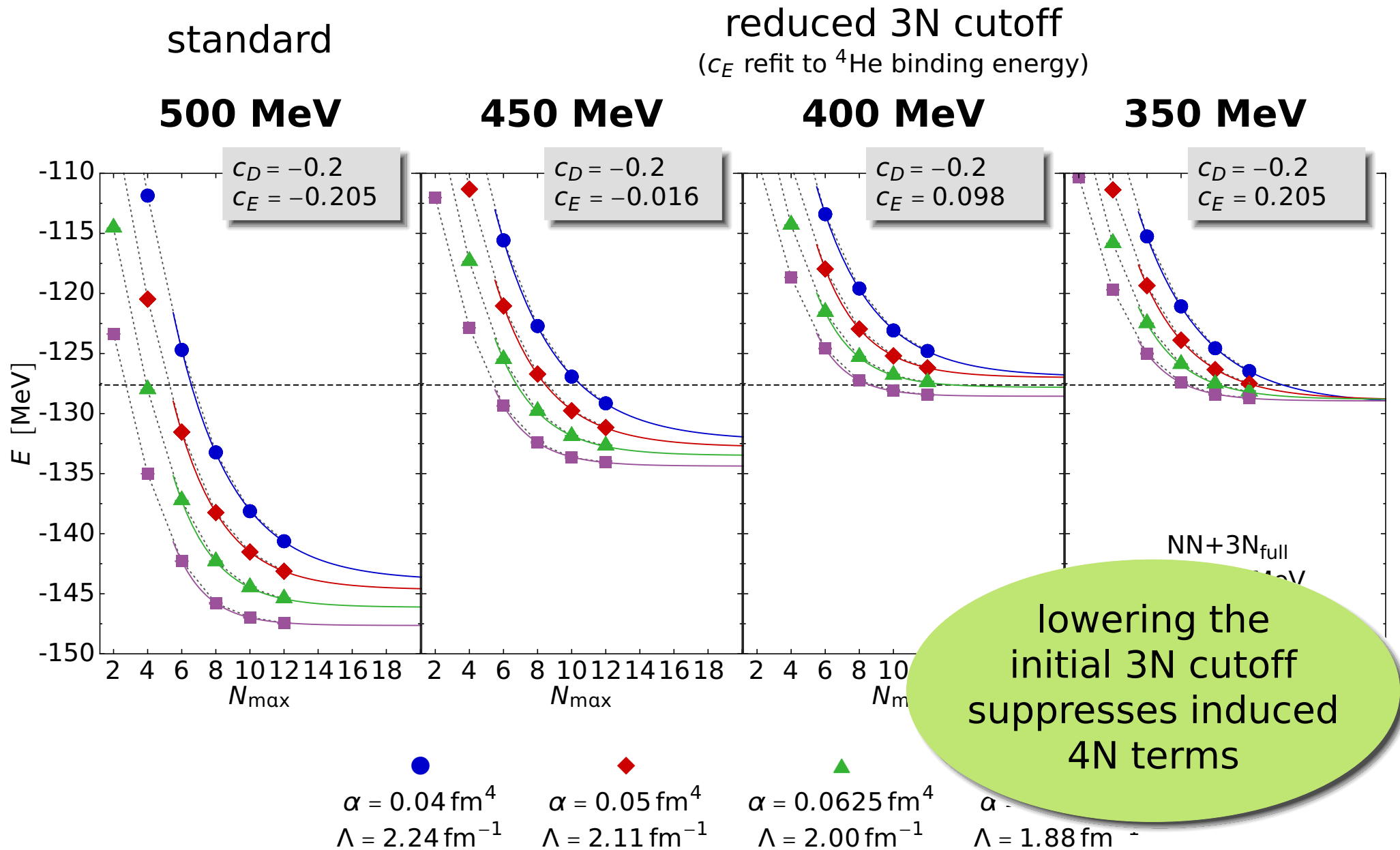


^{16}O : Ground-State Energies

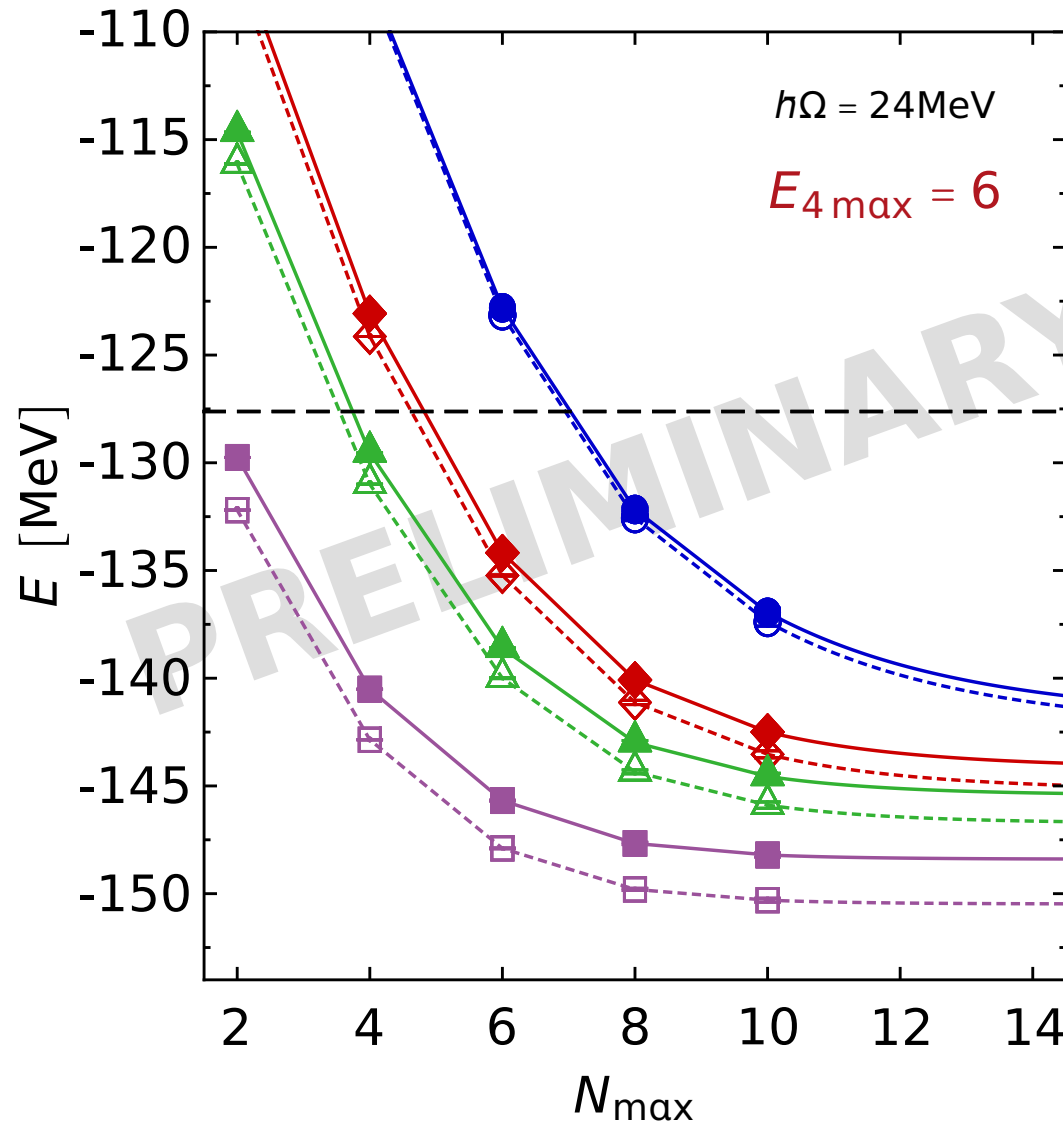
Roth, et al; PRL 107, 072501 (2011)



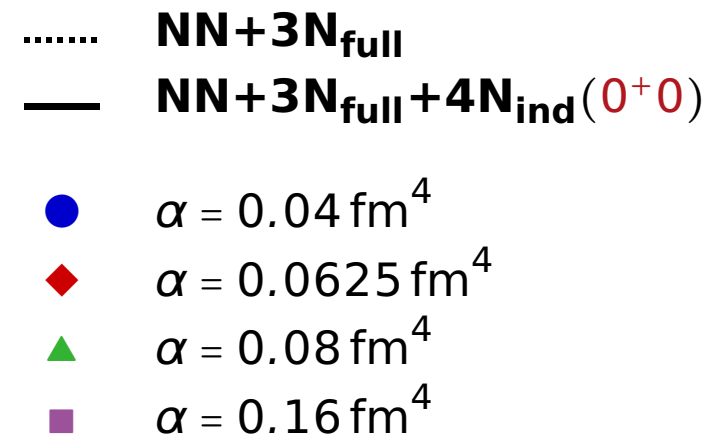
^{16}O : Lowering the Initial 3N Cutoff



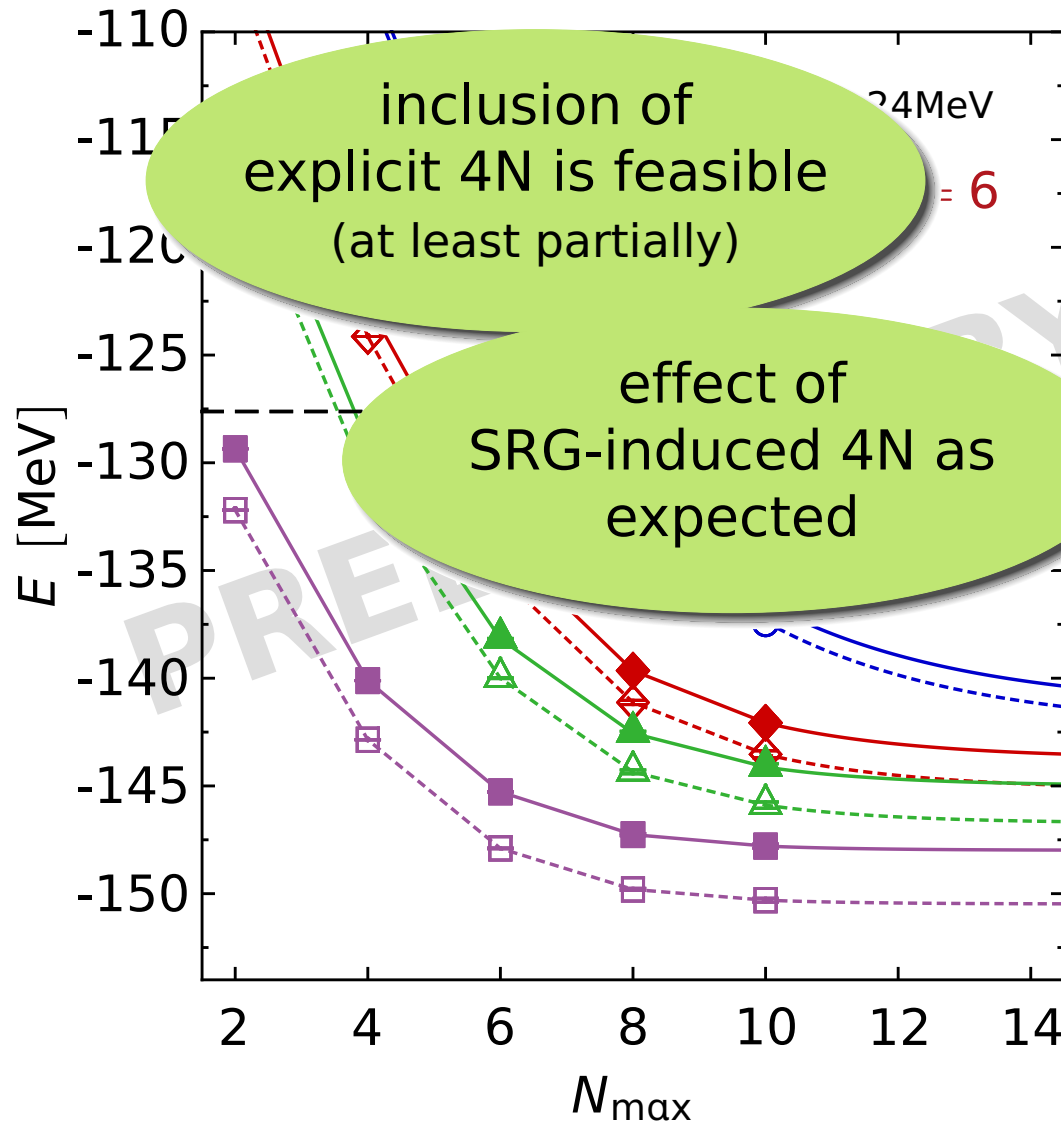
^{16}O : Explicit Inclusion of Induced 4N



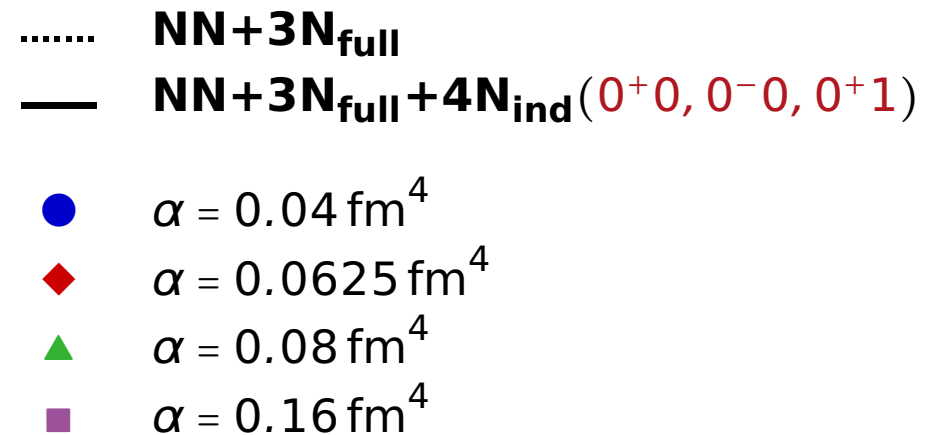
- induced 4N from SRG evolution in four-body Jacobi-HO
- transformation to m - or JT -scheme 4N matrix elements
- explicit 4N terms in IT-NCSM



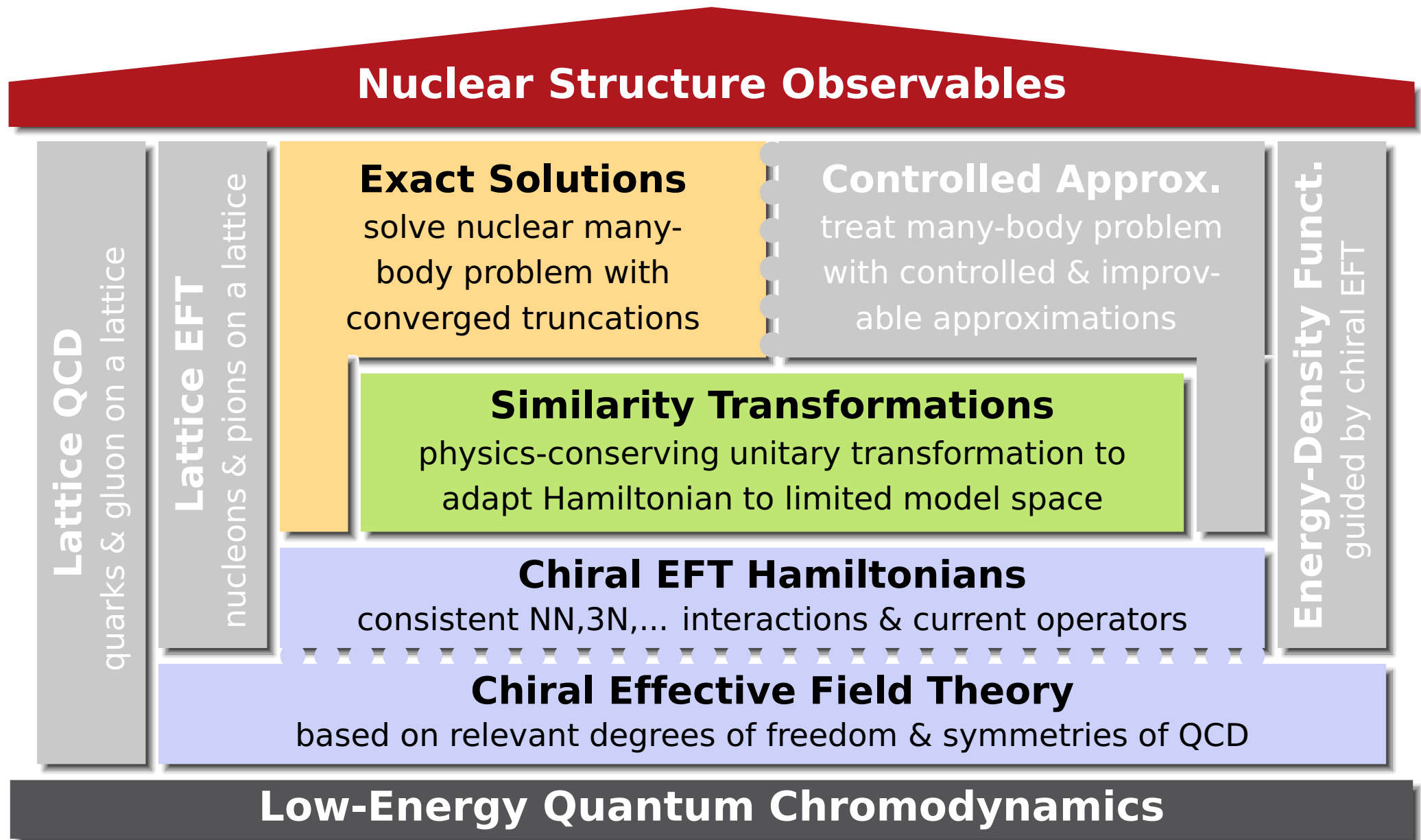
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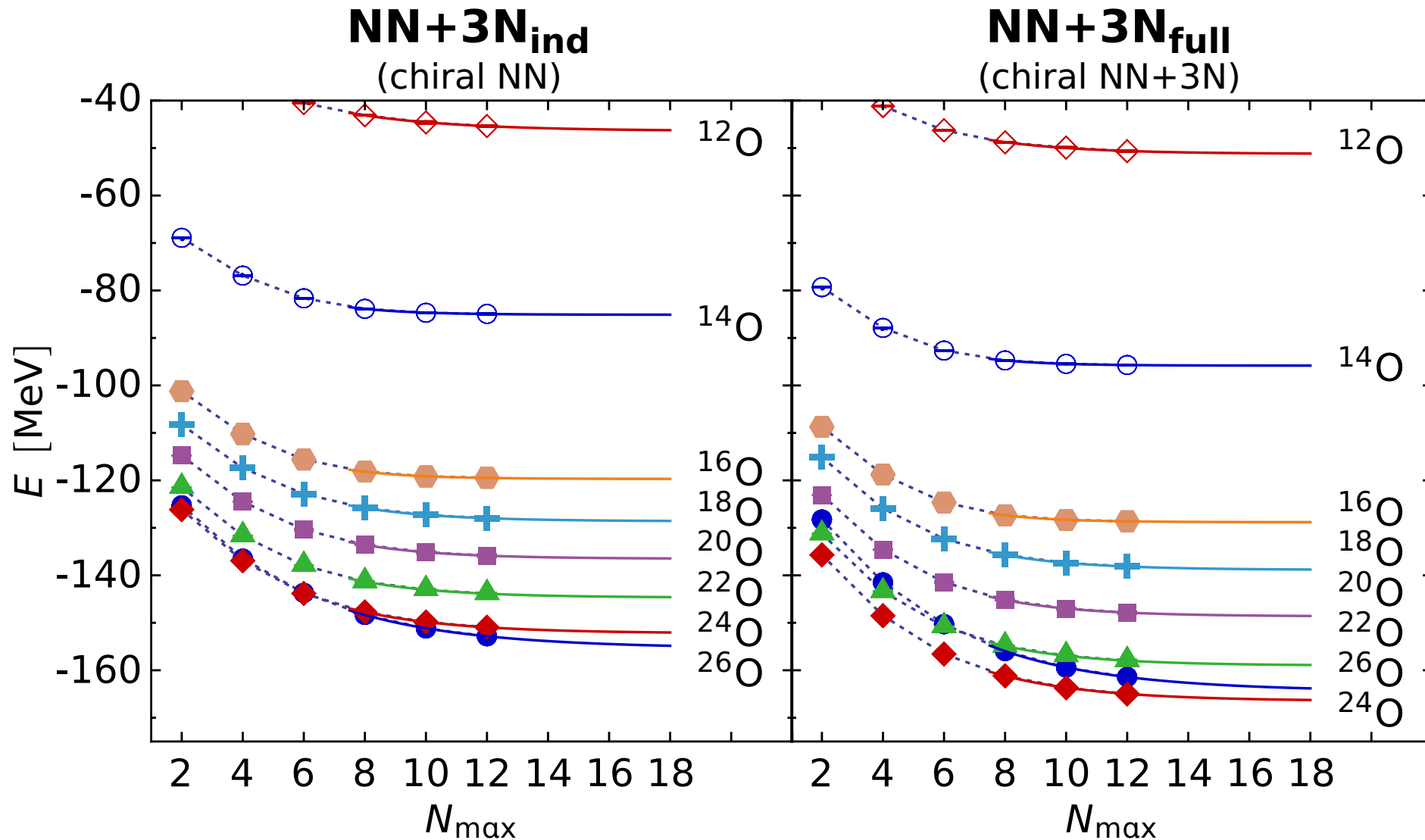


Ab Initio Nuclear Structure



Ground States of Oxygen Isotopes

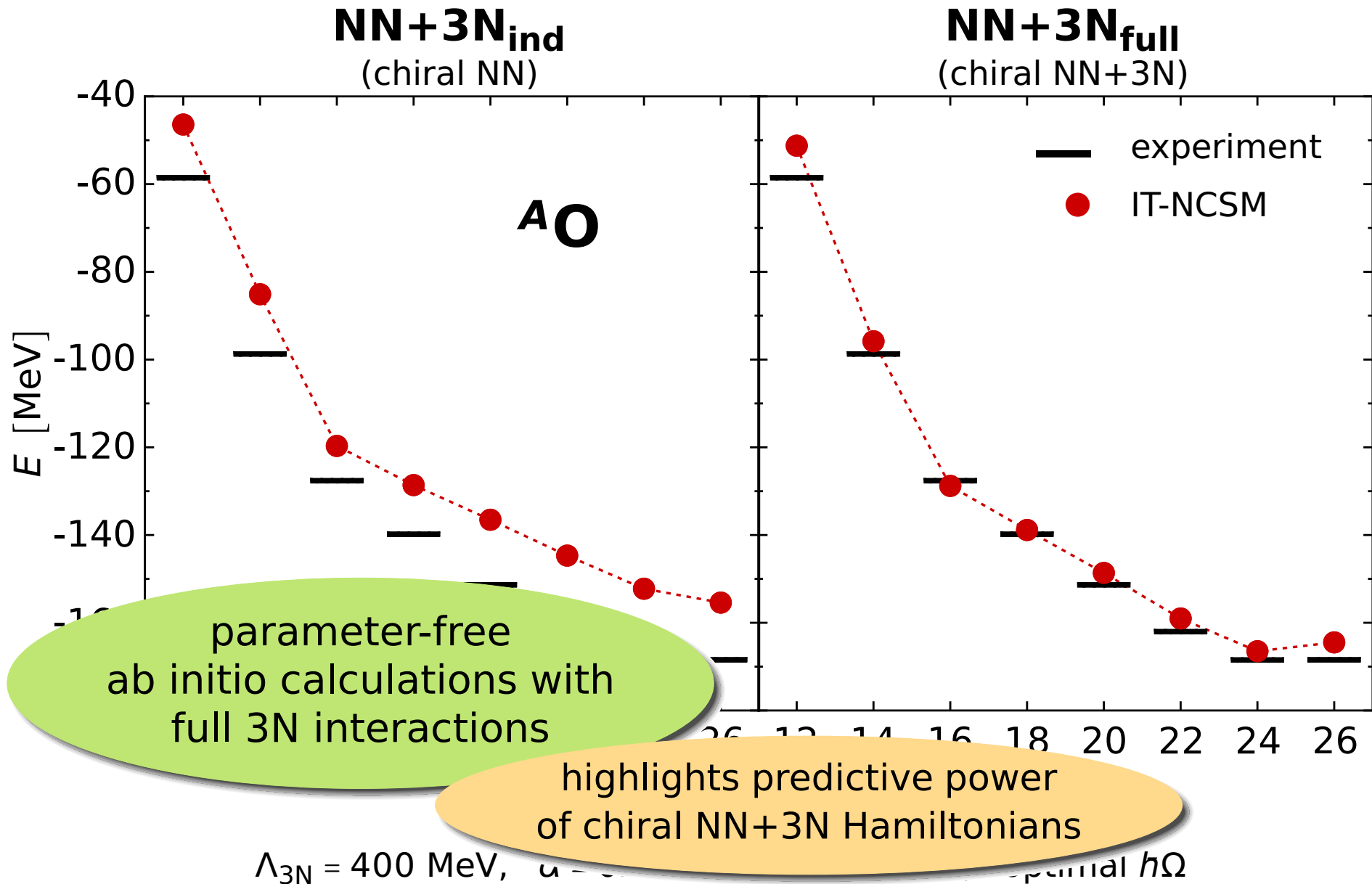
Hergert et al., PRL (2013); arXiv:1302.7294



$\Lambda_{3\text{N}} = 400 \text{ MeV}$, $\alpha = 0.08 \text{ fm}^4$, $E_{3\text{max}} = 14$, optimal $\hbar\Omega$

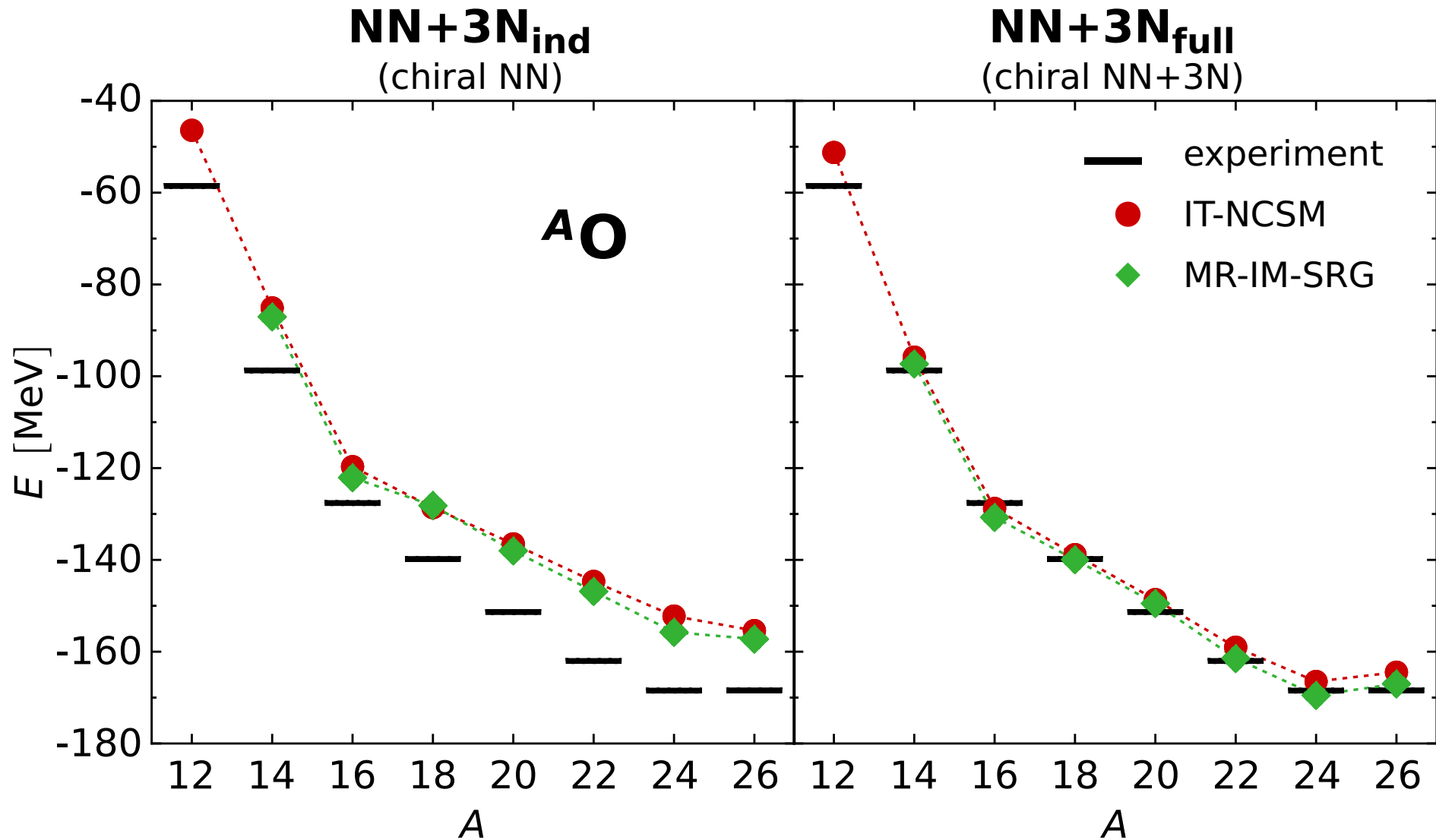
Ground States of Oxygen Isotopes

Hergert et al., PRL (2013); arXiv:1302.7294



Ground States of Oxygen Isotopes

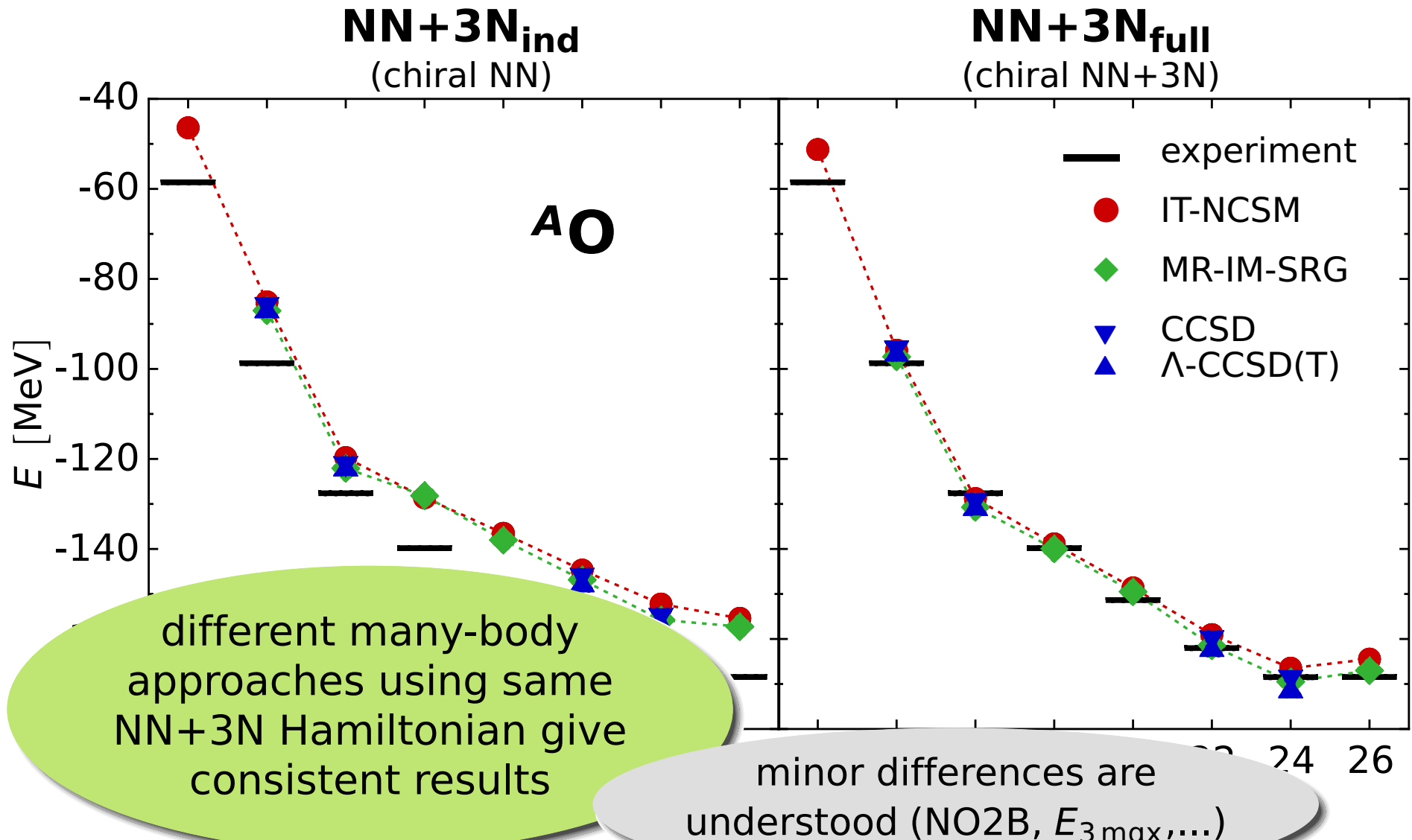
Hergert et al., PRL (2013); arXiv:1302.7294



$\Lambda_{3N} = 400$ MeV, $\alpha = 0.08 \text{ fm}^4$, $E_{3\text{max}} = 14$, optimal $h\Omega$

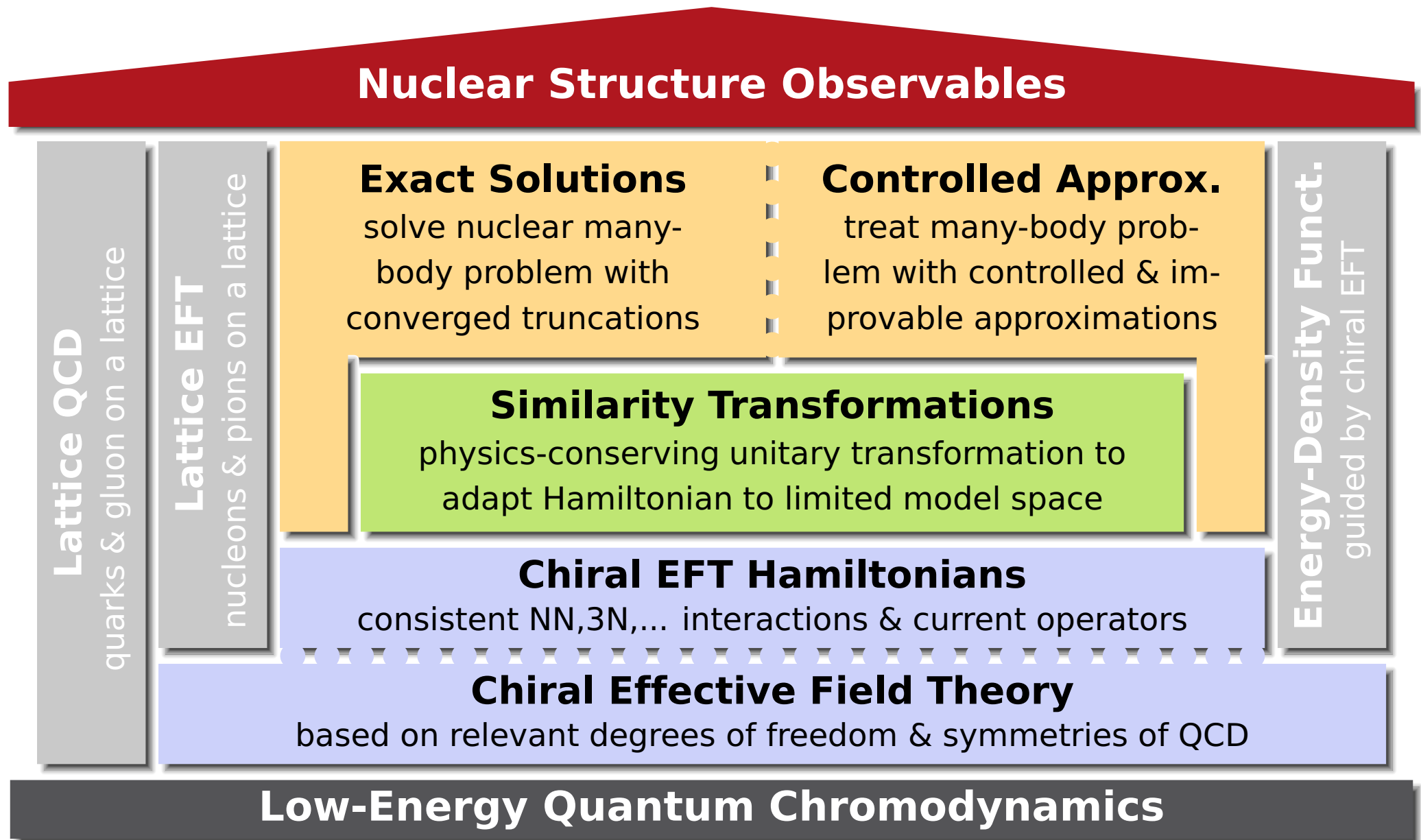
Ground States of Oxygen Isotopes

Hergert et al., PRL (2013); arXiv:1302.7294



$\Lambda_{3N} = 400$ MeV, $\alpha = 0.08$ fm⁻¹, $E_{3\text{max}} = 1\pi$, spatial $\hbar\Omega$

Ab Initio Nuclear Structure



Coupled-Cluster Method

Coester, Kuemmel, Bishop, Dean, Piecuch, Walet, Papenbrock, Hagen, Binder,...

CC is one of the most efficient methods for the description of ground states of medium-mass or heavy closed-shell nuclei

- many-body state parametrized as **exponential wave operator** applied to single-determinant **reference state** $|\Phi_{\text{ref}}\rangle$

$$|\Psi_{\text{CC}}\rangle = \Omega |\Phi_{\text{ref}}\rangle = \exp(T_1 + T_2 + T_3 + \dots + T_A) |\Phi_{\text{ref}}\rangle$$

- truncation with respect to n -particle- n -hole **excitation operators** T_n
- solve **non-linear system** of equations for the amplitudes in $T_1, T_2, T_3, \dots$
- extensions to near-closed-shell nuclei and excited states through **equations-of-motion methods**
- we have developed a parallelized CC code for **CCSD and Λ -CCSD(T)**

Inclusion of 3N Interactions

Roth, et al., PRL 109, 052501 (2012); Binder et al., PRC 87, 021303(R) (2013)

■ **premium option: explicit 3N**

- extend coupled-cluster equations for explicit 3N interactions
- CCSD-3B, Λ -CCSD(T)-3B are feasible, but much more expensive

■ **low-cost option: normal-ordered two-body approximation**

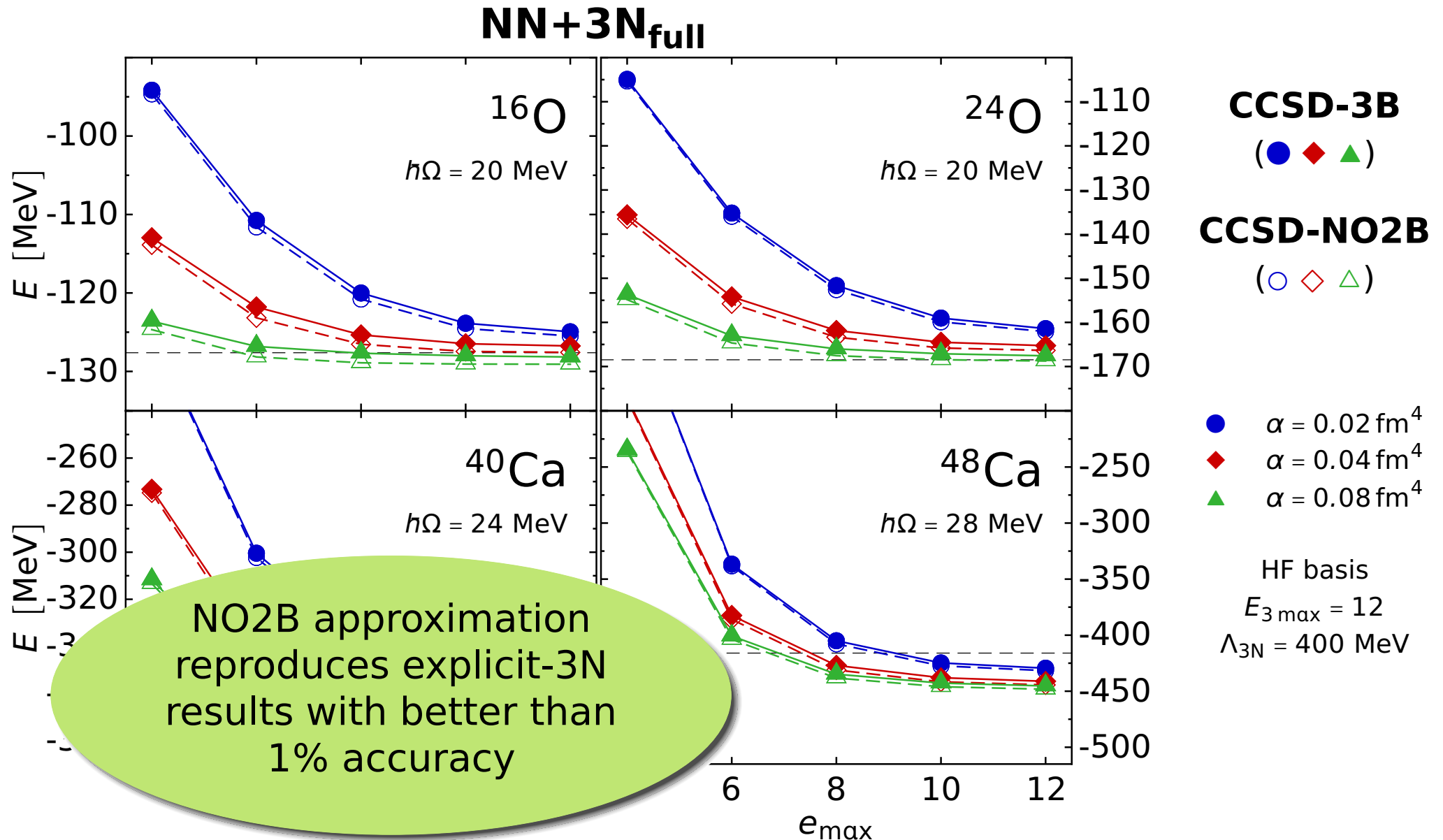
- write 3N interaction in normal-ordered form with respect to the actual A-body reference determinant (HF state)

$$\begin{aligned} V_{3N} &= \sum V_{\circ\circ\circ\circ\circ\circ}^{3N} a_{\circ}^{\dagger} a_{\circ}^{\dagger} a_{\circ}^{\dagger} a_{\circ} a_{\circ} a_{\circ} \\ &= W^{0B} + \sum W_{\circ\circ}^{1B} \{a_{\circ}^{\dagger} a_{\circ}\} + \sum W_{\circ\circ\circ}^{2B} \{a_{\circ}^{\dagger} a_{\circ}^{\dagger} a_{\circ} a_{\circ}\} \\ &\quad + \sum W_{\circ\circ\circ\circ\circ\circ}^{3B} \{a_{\circ}^{\dagger} a_{\circ}^{\dagger} a_{\circ}^{\dagger} a_{\circ} a_{\circ} a_{\circ}\} \end{aligned}$$

- discard normal-ordered three-body term and use two-body coupled-cluster formalism

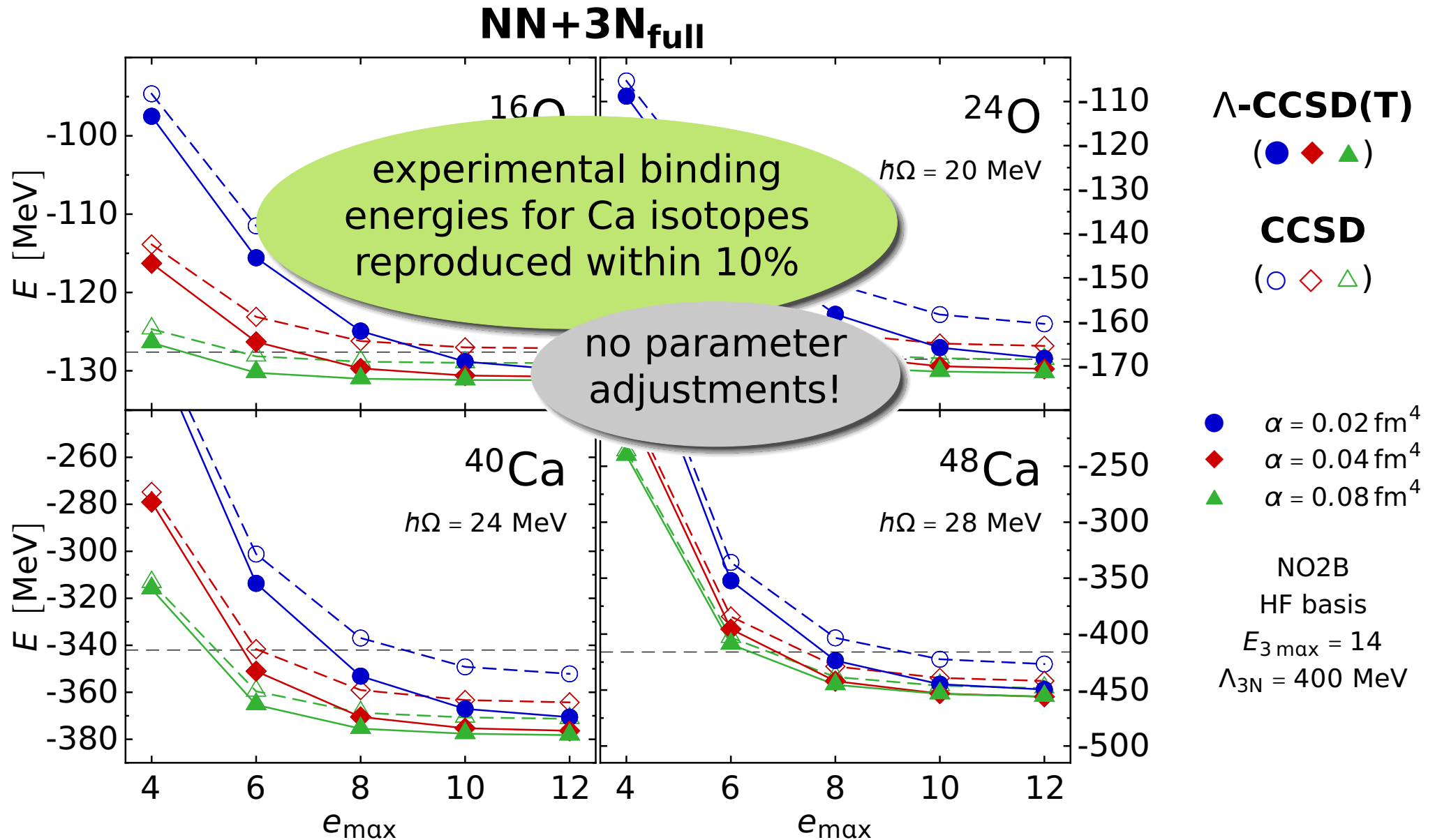
CCSD with Explicit 3N Interactions

Roth, et al., PRL 109, 052501 (2012); Binder et al., PRC 87, 021303(R) (2013)

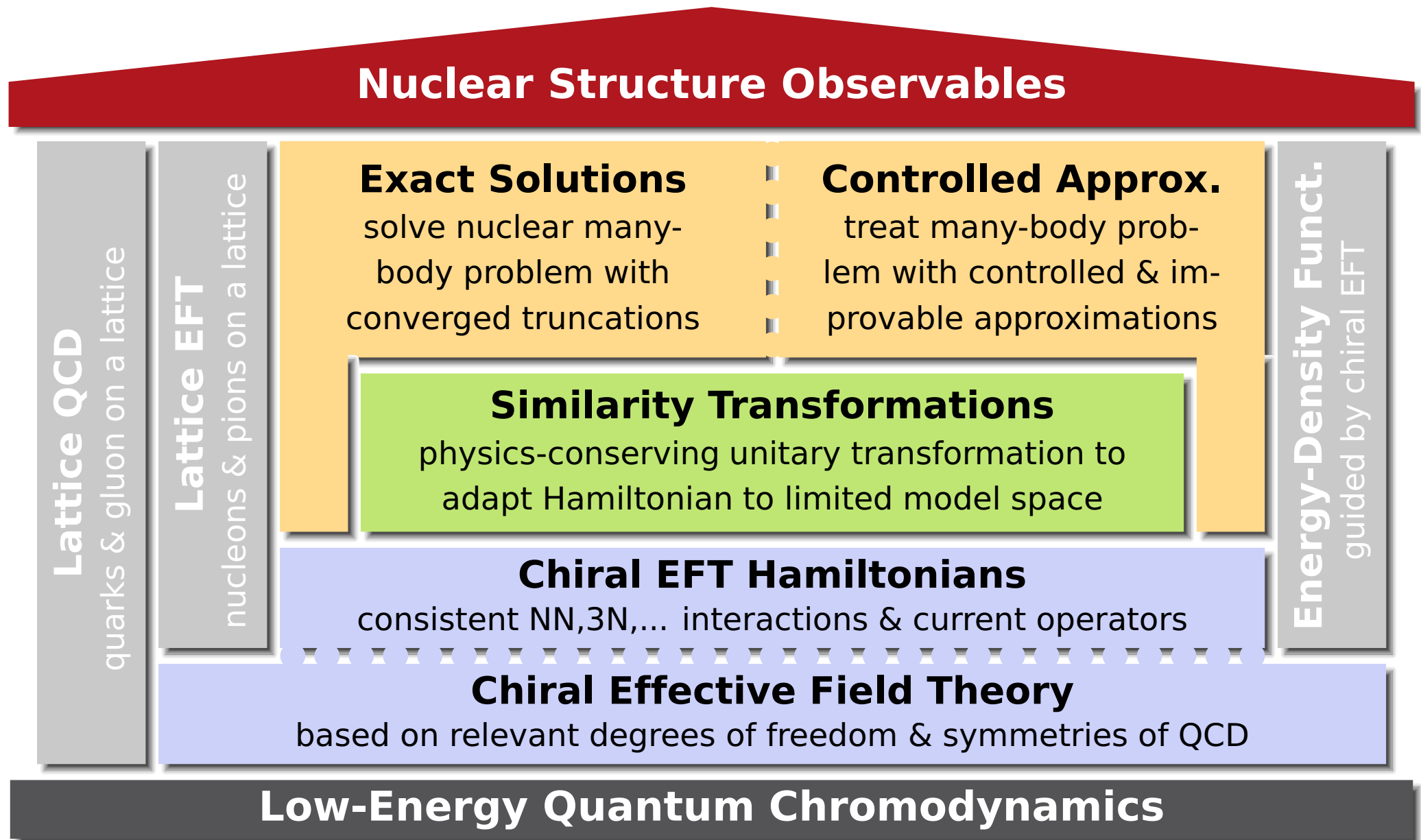


Λ -CCSD(T) with NO2B Approximation

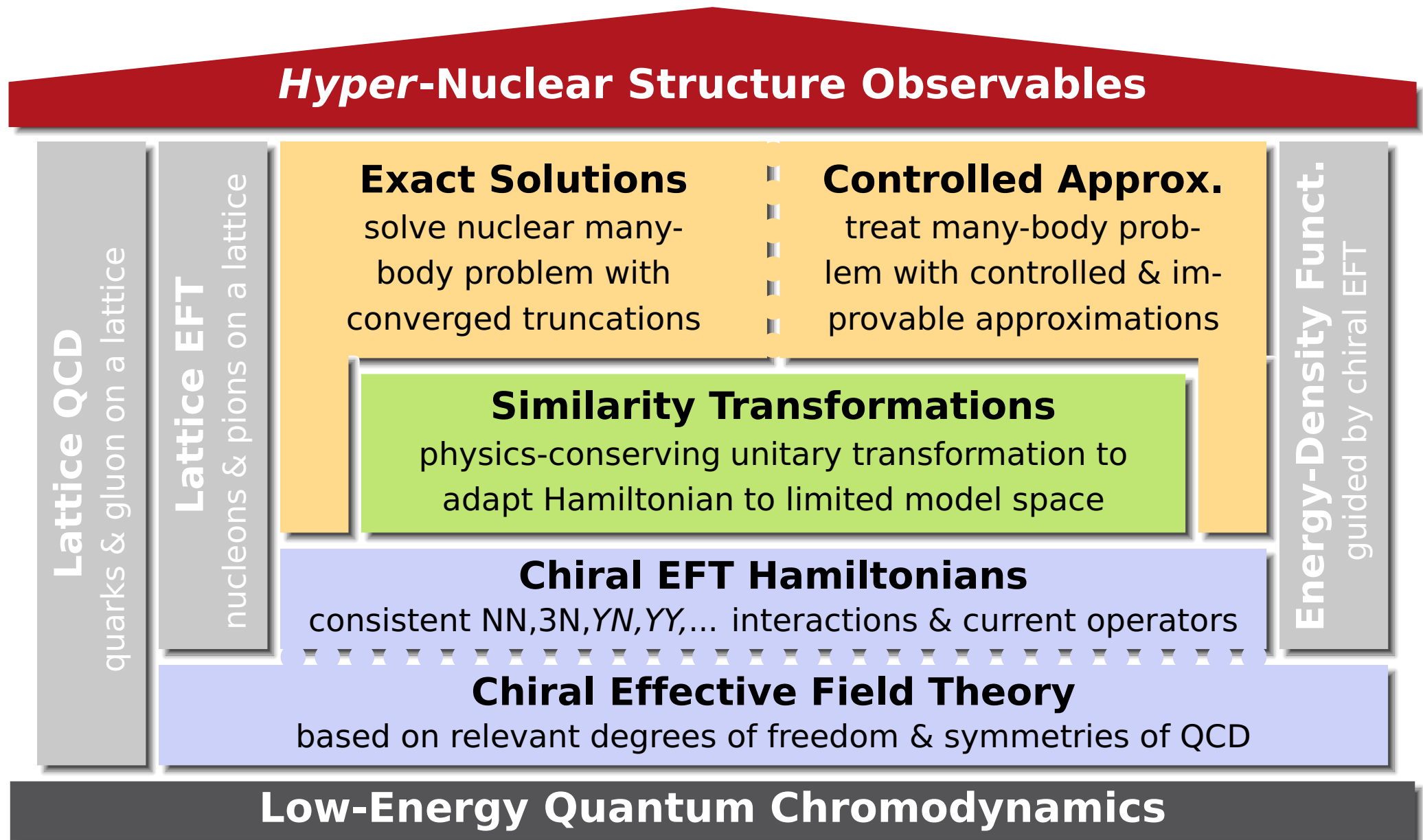
Roth, et al., PRL 109, 052501 (2012); Binder et al., PRC 87, 021303(R) (2013)



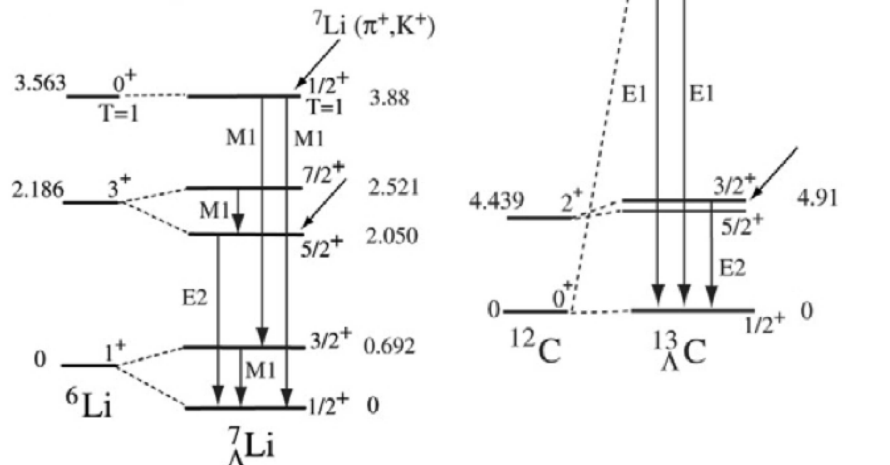
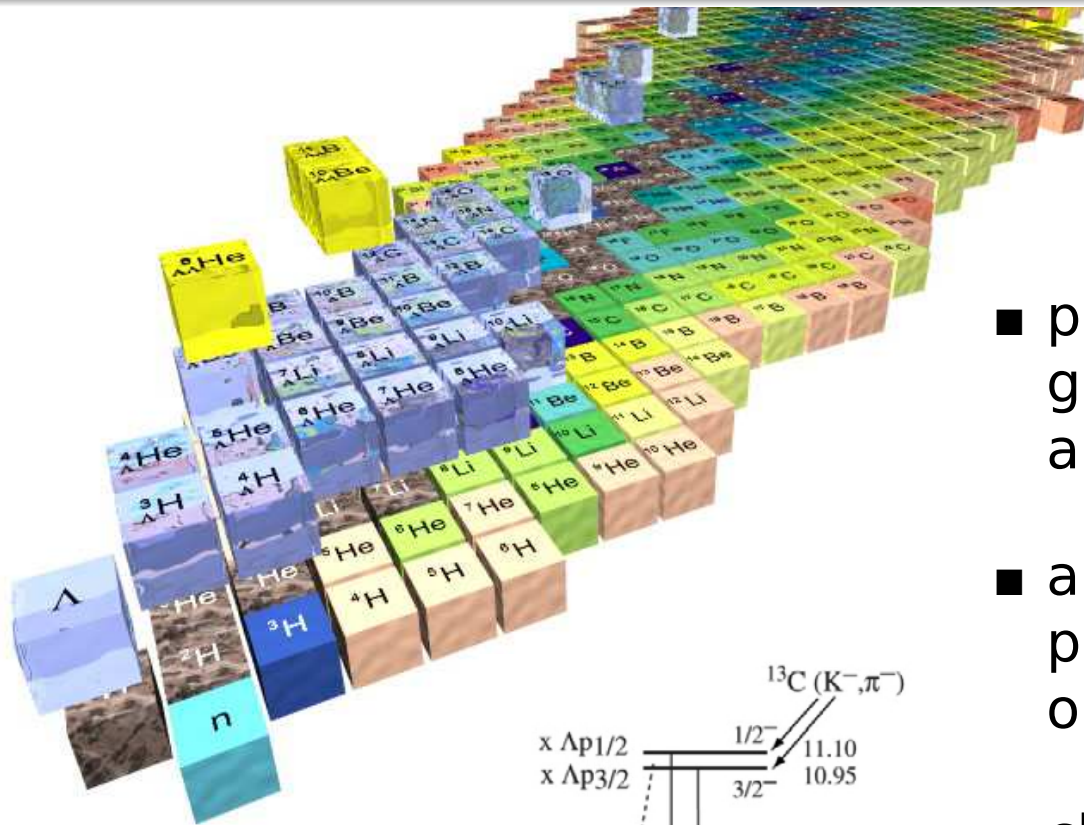
Ab Initio Nuclear Structure



Ab Initio *Hyper*-Nuclear Structure



Motivation: Hypernuclear Structure



- precision data on hypernuclear ground states and spectroscopy are available
- ab initio few-body ($A \lesssim 4$) and phenomenological shell model or cluster calculations so far
- chiral EFT interactions including hyperons are being constructed
- constrain YN & YY interaction by ab initio hypernuclear structure calculations

Ab Initio Toolbox

■ Hamiltonian from chiral EFT

- NN+3N: standard chiral Hamiltonian (Entem&Machleidt, Navrátil)
- YN: LO chiral interaction (Haidenbauer et al.), NLO in progress

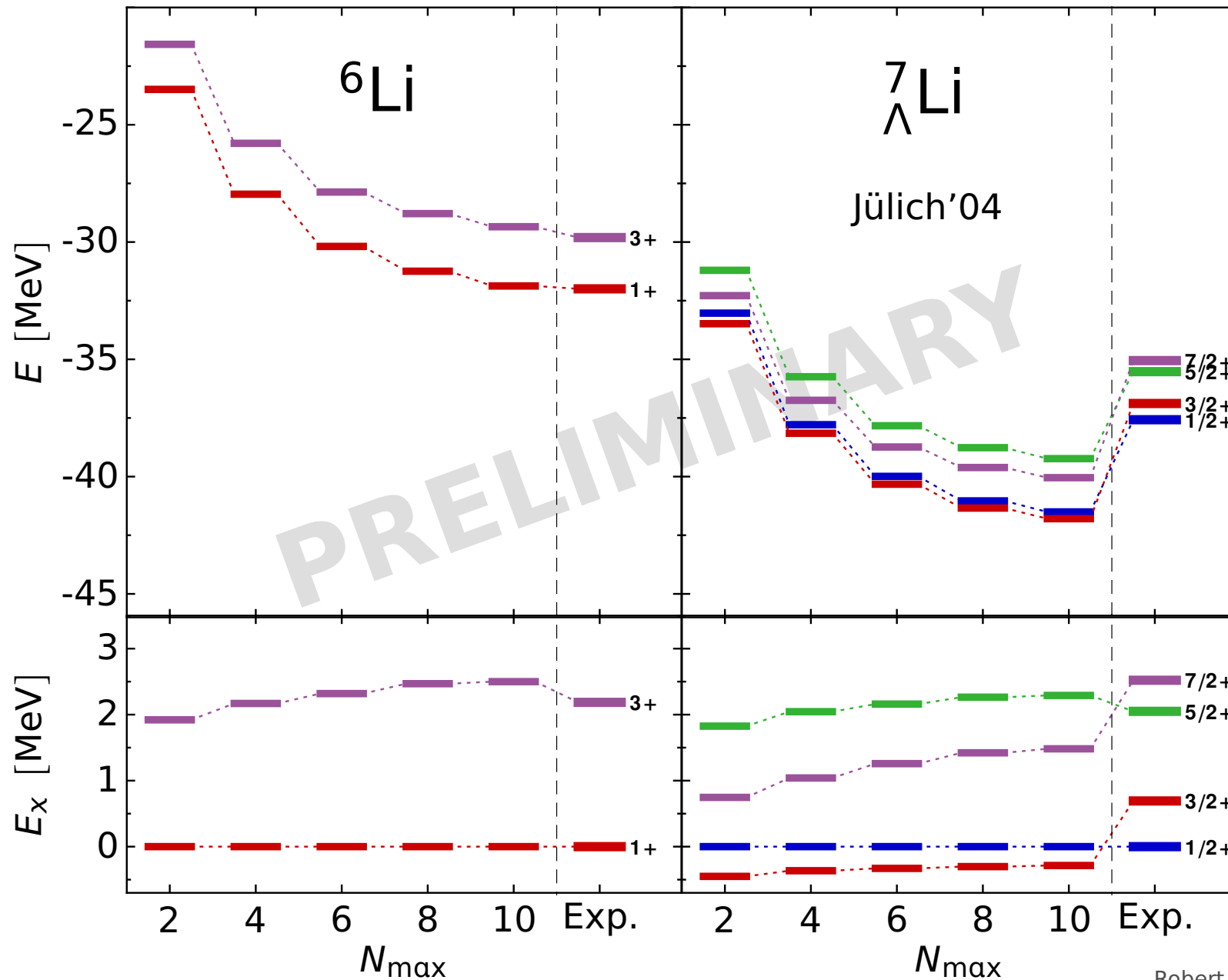
■ Similarity Renormalization Group

- consistent SRG-evolution of NN, 3N, YN interactions
- using particle basis and including $\Lambda\Sigma$ -coupling (larger matrices)
- Λ - Σ mass difference and $p\Sigma^\pm$ Coulomb included consistently

■ Importance Truncated No-Core Shell Model

- include explicit $(p, n, \Lambda, \Sigma^+, \Sigma^0, \Sigma^-)$ with physical masses
- larger model spaces easily tractable with importance truncation
- all p-shell single- Λ hypernuclei are accessible

Application: ${}^7_\Lambda\text{Li}$



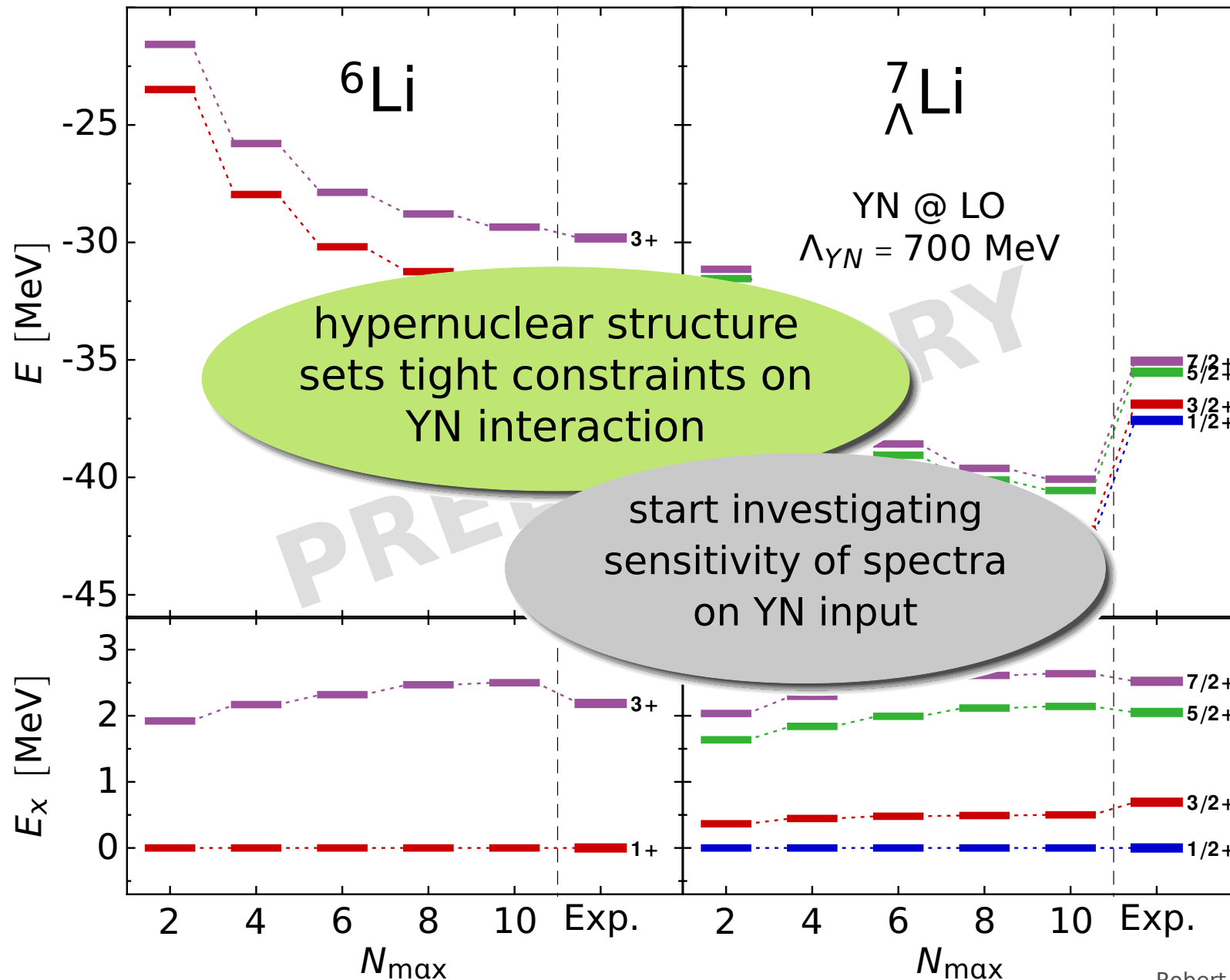
NN @ N3LO
Entem&Machleidt
 $\Lambda_{NN} = 500$ MeV

3N @ N2LO
Navratil
 $\Lambda_{3N} = 500$ MeV
triton fit

Jülich'04
Haidenbauer et al.
scatt. & hypertriton

$\alpha = 0.08 \text{ fm}^4$
 $\hbar\Omega = 20$ MeV

Application: ${}^7_\Lambda\text{Li}$



NN @ N3LO
Entem&Machleidt
 $\Lambda_{NN} = 500 \text{ MeV}$

3N @ N2LO
Navratil
 $\Lambda_{3N} = 500 \text{ MeV}$
triton fit

YN @ LO
Haidenbauer et al.
 $\Lambda_{YN} = 700 \text{ MeV}$
scatt. & hypertriton

$\alpha = 0.08 \text{ fm}^4$
 $\hbar\Omega = 20 \text{ MeV}$

Conclusions

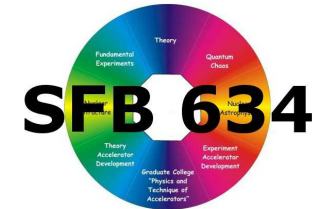
Conclusions

- new era of **ab-initio nuclear structure and reaction theory** connected to QCD via chiral EFT
 - chiral EFT as universal starting point... propagate uncertainties & provide feedback
- consistent **inclusion of 3N & 4N interactions** in similarity transformations & many-body calculations
 - breakthrough in treatment of 3N & 4N matrix elements
- **innovations in many-body theory**: extended reach of exact methods & improved control over approximations
 - versatile toolbox for different observables & mass ranges
- many **exciting applications** ahead...

Epilogue

■ thanks to my group & my collaborators

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LLNL Livermore, USA
- P. Piecuch
Michigan State University, USA
- **H. Hergert**, K. Hebeler
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- C. Forssén
Chalmers University, Sweden
- H. Feldmeier, T. Neff
GSI Helmholtzzentrum



Deutsche
Forschungsgemeinschaft
DFG



COMPUTING TIME



Happy Birthday James !!!

...and many more happy & productive visits to Darmstadt

