POSITRONIUM IN BASIS LIGHT FRONT QUANTIZATION

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INTRODUCTION

Light Front Bound State Equation

$$P^{2} |\Psi\rangle = M^{2} |\Psi\rangle \qquad H \equiv P^{2}$$

Fock Sector Decomposition of State Vector

$$\left|e^{-}e^{+}\right\rangle_{phys} = a\left|e^{-}e^{+}\right\rangle + b\left|e^{-}e^{+}\right\rangle + c\left|e^{-}e^{+}\right\rangle + d\left|\gamma\right\rangle + e\left|e^{-}e^{+}e^{-}e^{+}\right\rangle + \cdots$$

- In general: Large, sparse matrix diagonalization problem
- Apply suitably adapted methods from ab initio nuclear structure
- Test Case: Positronium

BASIS LIGHT FRONT QUANTIZATION

Transverse Harmonic Oscillator Basis

- Eigenfunctions of 2D potential $V = \frac{1}{2}M\Omega^2 \mathbf{r}^2$
- Every value of "oscillator energy" Ω determines a complete basis
 - Energy scale $b = \sqrt{M\Omega}$

• Restrict number of oscillator quanta $\sum (2n_i + |m_i| + 1) \le N_{\max}$



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BASIS LIGHT FRONT QUANTIZATION

- Discretize longitudinal momentum $p^+ = \frac{2\pi}{L}j$
 - Longitudinal box of length L
 - (anti)periodic boundary conditions
- Select value of total longitudinal momentum
 - Parameterize by dimensionless integer $P^+ = \frac{2\pi}{T}K$
 - K determines resolution of longitudinal momentum grid
- Pick total angular momentum projection "M-scheme"

$$\sum_{i} (m_i + s_i) = M_J$$

- For QCD, choose color singlet configurations
- Single-Particle Coordinates
 - Use specially chosen coordinates to guarantee exact CM factorization in truncated basis \mathbf{p}_i

$$\mathbf{q}_i = \frac{\mathbf{p}_i}{\sqrt{x_i}} \qquad \qquad \mathbf{s}_i = \sqrt{x_i} \mathbf{r}_i$$

$$x_i = \frac{p_i^+}{P^+} = \frac{j_i}{K}$$

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HAMILTONIAN MATRIX



EFFECTIVE POTENTIAL

■ Projection operators P and Q
$$|e^{-}e^{+}\rangle$$
 $|e^{-}e^{+}\rangle$
 $P^{2}|\Psi_{i}\rangle = M_{i}^{2}|\Psi_{i}\rangle$ $|e^{-}e^{+}\rangle$ PHP PHQ
 $H \equiv P^{2}$ $PHP \equiv H_{00}$
 $PHQ \equiv H_{01}$
 $QHP \equiv H_{10}$ $|e^{-}e^{+}\gamma\rangle$ QHP QHQ
 $QHQ \equiv H_{11}$

Eigenvalue equation can be rewritten:

 $H_{\text{eff}}(\omega) |\psi_i(\omega)\rangle_0 = \widetilde{M}_i^2(\omega) |\psi_i(\omega)\rangle_0$

$$H_{\text{eff}}(\omega) \equiv H_{00} + H_{01} \frac{1}{\omega - H_{11}} H_{10}$$

$$\omega = M_i^2$$

EFFECTIVE POTENTIAL

$$H_{\rm eff}(\omega) \equiv H_{00} + H_{01} \frac{1}{\omega - H_{11}} H_{10}$$



In special case where H₁₁ is diagonal

$$\left\langle f \left| H_{eff} \left(\omega \right) \right| i \right\rangle = \left\langle f \left| H_{00} \right| i \right\rangle + \sum_{|n\rangle} \frac{\left\langle f \left| H_{01} \right| n \right\rangle \left\langle n \left| H_{10} \right| i \right\rangle}{\omega - \left\langle n \left| H_{11} \right| n \right\rangle}$$

- **Resolvent operator** $(\omega H_{11})^{-1}$ not diagonal in H.O. basis
 - Work in momentum space initially

SMALL-x DIVERGENCES

Perturbative calculation of scattering amplitude



- Each term has a divergence at small x
 - Divergences cancel in perturbative calculation
 - Result is identical to scattering amplitude obtained via usual Feynman rules in equal-time quantization

TWO-BODY EFFECTIVE INTERACTION

• Consider just second term of effective interaction $\langle f | H_{eff}^{(2)}(\omega) | i \rangle = \sum_{|n\rangle} \frac{\langle f | H_{01} | n \rangle \langle n | H_{10} | i \rangle}{\omega - E_n}$

Sum reduces to sum over polarization states of photon

$$\sum_{\lambda_g} \varepsilon_{\mu}(k_g, \lambda_g) \varepsilon_{\nu}^*(k_g, \lambda_g) = -g_{\mu\nu} + (k_{g,\mu}\eta_{\nu} + k_{g,\nu}\eta_{\mu})/k_g^{\kappa}\eta_{\kappa}$$

$$\eta^{\mu} = (\eta^{\mu}, \eta_{\perp}, \eta^{\mu}) = (0, \mathbf{0}_{\perp}, 2)$$

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Expected cancellation is achieved only if:

$$\omega = \frac{E_i + E_f}{2}$$

EXPRESSION FOR EFFECTIVE POTENTIAL

Final result

$$\left\langle f \left| H_{eff}^{(2)} \left(\omega = \frac{E_i + E_f}{2} \right) \right| i \right\rangle = \alpha \frac{\delta_{x_1 + x_2}^{x_1 + x_2}}{K} \int \frac{d^2 q_1}{(2\pi)^2} \frac{d^2 q_2}{(2\pi)^2} \frac{d^2 q_1'}{(2\pi)^2} \frac{d^2 q_2'}{(2\pi)^2} (2\pi)^2 \delta^{(2)} \left(\sqrt{x_1} q_1 + \sqrt{x_2} q_2 - \sqrt{x_1'} q_1' + \sqrt{x_2'} q_2' \right) \right. \\ \left. \left. \frac{\Psi_{n_1}^{m_1} \left(q_1 \right) \Psi_{n_2}^{m_2} \left(q_2 \right) \Psi_{n_1'}^{*m_1'} \left(q_1' \right) \Psi_{n_2'}^{*m_2'} \left(q_2' \right) \overline{u} \left(1' \right) \gamma^{\mu} u \left(1 \right) \overline{v} \left(2 \right) \gamma_{\mu} v \left(2' \right) \right. \\ \left. \frac{\left(\frac{x_1 - x_1'}{2} \right) \left[\left(\frac{x_1 q_1^2 + m^2}{x_1} + \frac{x_1' q_1'^2 + m^2}{x_1'} - \frac{\left(\sqrt{x_1} q_1 - \sqrt{x_1'} q_1' \right)^2 + \mu^2}{x_1 - x_1'} \right) - \left(\frac{x_2 q_2^2 + m^2}{x_2} + \frac{x_2' q_2'^2 + m^2}{x_2'} - \frac{\left(\sqrt{x_2} q_2 - \sqrt{x_2'} q_2' \right)^2 + \mu^2}{x_2 - x_2'} \right) \right]$$

$$I = -\frac{1}{x_1 + x_2} \cdot \frac{1}{2\pi} \cdot \delta_{m_1 + m_2}^{m_1' + m_2'} \cdot \sum (TM)(TM)(TM) \int dx dy \frac{e^{-(x+y)/2}L(x)L(y)}{(\alpha - \beta)^2 x + (\alpha + \beta)^2 y + \frac{\Lambda}{b^2} + \varepsilon}$$

$$\alpha = \sqrt{\frac{x_1' x_2}{x_1 + x_2}} \qquad \Delta = m^2 (x_1 - x_1')^2 \Big[\frac{1}{x_1' x_1} + \frac{1}{x_2' x_2}\Big] \ge 0 \qquad \varepsilon = 2\frac{\mu^2}{b^2}$$

- 3-fold summation
- 2D Numerical Integral
- Infrared Cutoff

BLFQ RESULTS FOR POSITRONIUM

 $\alpha = 0.3$ $\mu = 0.1 \text{m}_f$ $b = 0.5 \text{m}_f$

$$\varepsilon = 2\frac{\mu^2}{b^2} = 0.08$$



$$E = a + be^{-c\sqrt{K}}$$

$\epsilon \rightarrow 0 LIMIT$



COMPARISON



GROUND STATE DISTRIBUTION FUNCTION



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SUMMARY

- Effective two-body photon exchange interaction
 - Cancellation of Light Front small-x divergences obtained only for a particular value of arbitrary parameter ω

Finite result in continuum limit

- Results for binding energy and hyperfine splitting of ground state are the same order of magnitude as non-relativistic QM predictions
- Overbinding for large α also found by H.C. Pauli's group (DLCQ)
- Future Work:
 - Heavy Quarkonia
 - Explicit dynamical photons in basis
 - Non-perturbative renormalization

HARMONIC OSCILLATOR BASIS

Harmonic Oscillator Basis for Coulomb Problem

• Wrong asymptotic behavior $e^{-\alpha r^2}$ vs. $e^{-\beta r}$



VARIATIONAL CALCULATION

■ ϵ → 0 limit is safe provided K → ∞

 $\alpha = 0.3$ 2.1 $\mu = 0.005 m_f$ 0 • K -> Infinity $N_{\rm max} = 6$ Ground State Energy (fermion mass units) K=13 0 0 2 0 0 0 0 0 0 0 0 0 ο 1.9 $\varepsilon = 2\frac{\mu^2}{b^2}$ 0 0.5 0 1.5 1 b

EXTRA

Neglect instantaneous interactions when corresponding dynamical exchange is not present in model space

